

Quantum State Transformation II

State Transition in Multi-Qubit System

- We start with a 2-qubit systems in **standard** or **computational basis**.
- The basis of such system is the tensor product of two 1-qubit bases -

$$\begin{aligned} & \{|0\rangle, |1\rangle\} \otimes \{|0\rangle, |1\rangle\} \\ = & \{|0\rangle \otimes |0\rangle, |0\rangle \otimes |1\rangle, |1\rangle \otimes |0\rangle, |1\rangle \otimes |1\rangle\} \\ = & \{|0\rangle |0\rangle, |0\rangle |1\rangle, |1\rangle |0\rangle, |1\rangle |1\rangle\} \\ = & \{|00\rangle, |01\rangle, |10\rangle, |11\rangle\} \end{aligned}$$

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2-Qubit Representation

- If we represent $|0\rangle = \begin{bmatrix} 1 \\ 0 \end{bmatrix}$ and $|1\rangle = \begin{bmatrix} 0 \\ 1 \end{bmatrix}$,

$$|00\rangle = \begin{bmatrix} 1 \\ 0 \end{bmatrix} \otimes \begin{bmatrix} 1 \\ 0 \end{bmatrix} = \begin{bmatrix} 1 & 1 \\ 0 & 0 \end{bmatrix} = \begin{bmatrix} 1 \\ 0 \\ 0 \\ 0 \end{bmatrix}$$

2-Qubit Representation

- Similarly,

$$|01\rangle = \begin{bmatrix} 0 \\ 1 \\ 0 \\ 0 \end{bmatrix}, \quad |10\rangle = \begin{bmatrix} 0 \\ 0 \\ 1 \\ 0 \end{bmatrix}, \quad |11\rangle = \begin{bmatrix} 0 \\ 0 \\ 0 \\ 1 \end{bmatrix}.$$

2-Qubit Unitary Transformation

- One may think of a **2-qubit unitary transformation** as a **tensor product** of two **1-qubit unitary transformations**.
- As an example consider

$$X \otimes Y = \begin{bmatrix} 0 & 1 \\ 1 & 0 \end{bmatrix} \otimes \begin{bmatrix} 0 & -i \\ i & 0 \end{bmatrix} = \begin{bmatrix} 0 & \begin{bmatrix} 0 & -i \\ i & 0 \end{bmatrix} \\ 1 & \begin{bmatrix} 0 & -i \\ i & 0 \end{bmatrix} \end{bmatrix} = \begin{bmatrix} 0 & -i & 0 & 0 \\ 0 & 0 & -i & 0 \\ 1 & 0 & 0 & -i \\ 0 & i & 0 & 0 \end{bmatrix}$$

2-Qubit Unitary Transformation

$$X \otimes Y = \begin{bmatrix} 0 & 0 & 0 & -i \\ 0 & 0 & i & 0 \\ 0 & -i & 0 & 0 \\ i & 0 & 0 & 0 \end{bmatrix}$$

The effect of $X \otimes Y$ on a two qubit state $|pq\rangle$ is same as $|(Xp) \otimes (Yq)\rangle$.

We consider an example where $p = \begin{bmatrix} 0 \\ 1 \end{bmatrix}$ and $q = \begin{bmatrix} 1 \\ 0 \end{bmatrix}$.

2-Qubit Unitary Transformation

$$\begin{aligned}
 (X \otimes Y) |10\rangle &= \begin{bmatrix} 0 & 1 \\ 1 & 0 \end{bmatrix} \otimes \begin{bmatrix} 0 & -i \\ i & 0 \end{bmatrix} \left(\begin{bmatrix} 0 \\ 1 \end{bmatrix} \otimes \begin{bmatrix} 1 \\ 0 \end{bmatrix} \right) \\
 &= \begin{bmatrix} 0 & 1 \\ 1 & 0 \end{bmatrix} \begin{bmatrix} 0 \\ 1 \end{bmatrix} \otimes \begin{bmatrix} 0 & -i \\ i & 0 \end{bmatrix} \begin{bmatrix} 1 \\ 0 \end{bmatrix} = \begin{bmatrix} 1 \\ 0 \end{bmatrix} \otimes \begin{bmatrix} 0 \\ i \end{bmatrix} \\
 &= \begin{bmatrix} 0 \\ i \\ 0 \\ 0 \end{bmatrix} = i |01\rangle = |0\rangle \otimes i |1\rangle
 \end{aligned}$$

Negation of the 1st qubit and Y on the second qubit.
This is same as -

2-Qubit Unitary Transformation

$$\begin{aligned}(X \otimes Y) |10\rangle &= \begin{bmatrix} 0 & 0 & 0 & -i \\ 0 & 0 & i & 0 \\ 0 & -i & 0 & 0 \\ i & 0 & 0 & 0 \end{bmatrix} \begin{bmatrix} 0 \\ 0 \\ 1 \\ 0 \end{bmatrix} \\ &= \begin{bmatrix} 0 \\ i \\ 0 \\ 0 \end{bmatrix}.\end{aligned}$$

2-Qubit Unitary Transformation

In general, if $p = \begin{bmatrix} a \\ b \end{bmatrix}$ and $q = \begin{bmatrix} c \\ d \end{bmatrix}$,

$$\begin{aligned}
 (X \otimes Y) |pq\rangle &= \begin{bmatrix} 0 & 0 & 0 & -i \\ 0 & 0 & i & 0 \\ 0 & -i & 0 & 0 \\ i & 0 & 0 & 0 \end{bmatrix} \begin{bmatrix} ac \\ ad \\ bc \\ bd \end{bmatrix} = \begin{bmatrix} -ibd \\ ibc \\ -iad \\ iac \end{bmatrix} \\
 &= \begin{bmatrix} b \\ a \end{bmatrix} \otimes \begin{bmatrix} -id \\ ic \end{bmatrix} = (Xp) \otimes (Yq).
 \end{aligned}$$

Note

- This type of transformations do not create any new dependency (entanglement) of qubits.
- But there are 2-qubit transformations that cannot be expressed as a tensor product of two 1-qubit transformations.

CNOT Gate

One of the most important of such transformations is **CNOT**. We have already shown (Boolean) that it cannot be expressed as a tensor product of two 1-qubit transformations.

$$\begin{bmatrix} a & b \\ c & d \end{bmatrix} \otimes \begin{bmatrix} p & q \\ r & s \end{bmatrix} = \begin{bmatrix} ap & aq & bp & bq \\ ar & as & br & bs \\ cp & cq & dp & dq \\ cr & cs & dr & ds \end{bmatrix} \neq \begin{bmatrix} 1 & 0 & 0 & 0 \\ 0 & 1 & 0 & 0 \\ 0 & 0 & 0 & 1 \\ 0 & 0 & 1 & 0 \end{bmatrix}$$

$dp = 0$ implies that either $d = 0$ or $p = 0$. If $d = 0$, then $dq = 0$ - not possible. If $p = 0$, then $ap = 0$ - also not possible; so a contradiction.

CNOT Gate

The **CNOT** transformation can be expressed as

$$\begin{aligned}
 |0\rangle\langle 0| \otimes I + |1\rangle\langle 1| \otimes X &= \begin{bmatrix} 1 & 0 \\ 0 & 0 \end{bmatrix} \otimes \begin{bmatrix} 1 & 0 \\ 0 & 1 \end{bmatrix} + \begin{bmatrix} 0 & 0 \\ 0 & 1 \end{bmatrix} \otimes \begin{bmatrix} 0 & 1 \\ 1 & 0 \end{bmatrix} \\
 &= \begin{bmatrix} 1 & 0 & 0 & 0 \\ 0 & 1 & 0 & 0 \\ 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 \end{bmatrix} + \begin{bmatrix} 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 1 \\ 0 & 0 & 1 & 0 \end{bmatrix} \\
 &= \begin{bmatrix} 1 & 0 & 0 & 0 \\ 0 & 1 & 0 & 0 \\ 0 & 0 & 0 & 1 \\ 0 & 0 & 1 & 0 \end{bmatrix}
 \end{aligned}$$

CNOT Gate

We have $p = \begin{bmatrix} a \\ b \end{bmatrix}$ and $q = \begin{bmatrix} c \\ d \end{bmatrix}$.

$$CNOT |pq\rangle = \begin{bmatrix} 1 & 0 & 0 & 0 \\ 0 & 1 & 0 & 0 \\ 0 & 0 & 0 & 1 \\ 0 & 0 & 1 & 0 \end{bmatrix} \begin{bmatrix} ac \\ ad \\ bc \\ bd \end{bmatrix} = \begin{bmatrix} ac \\ ad \\ bd \\ bc \end{bmatrix}.$$

$$\left(\begin{bmatrix} 1 & 0 & 0 & 0 \\ 0 & 1 & 0 & 0 \\ 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 \end{bmatrix} + \begin{bmatrix} 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 1 \\ 0 & 0 & 1 & 0 \end{bmatrix} \right) \begin{bmatrix} ac \\ ad \\ bc \\ bd \end{bmatrix} = \begin{bmatrix} ac \\ ad \\ 0 \\ 0 \end{bmatrix} + \begin{bmatrix} 0 \\ 0 \\ bd \\ bc \end{bmatrix} = \begin{bmatrix} ac \\ ad \\ bd \\ bc \end{bmatrix}$$

Note

- We have seen that the effect of $X \otimes Y$ on a pair of qubits is an application of X on the first qubit and an application of Y on the second qubit. One action does not influence the other.
- On the other hand in case of CNOT, the first qubit influences the action on the second qubit - it is either **identity** or **NOT**.

Note

- If $p = \begin{bmatrix} a \\ b \end{bmatrix} = \begin{bmatrix} 1 \\ 0 \end{bmatrix}$, the new state comes from $\begin{bmatrix} ac \\ ad \\ 0 \\ 0 \end{bmatrix}$, and no change in the order of c, d .
- If $p = \begin{bmatrix} a \\ b \end{bmatrix} = \begin{bmatrix} 0 \\ 1 \end{bmatrix}$, the new state comes from $\begin{bmatrix} 0 \\ 0 \\ bd \\ bc \end{bmatrix}$, and the order of c, d are reversed.

CNOT versus $X \otimes Y$

Consider the state $|x\rangle = \frac{1}{\sqrt{2}}(|00\rangle + |10\rangle)$.

- $(X \otimes Y) \left(\frac{1}{\sqrt{2}}(|00\rangle + |10\rangle) \right) = \frac{i}{\sqrt{2}}(|11\rangle + |01\rangle)$.
- $CNOT |x\rangle = CNOT \left(\frac{1}{\sqrt{2}}(|00\rangle + |10\rangle) \right) = \frac{1}{\sqrt{2}}(|00\rangle + |11\rangle)$.

The initial state was **not entangled**. The state after the transformation $X \otimes Y$ is also **not entangled**, but CNOT creates an **entangled state**.

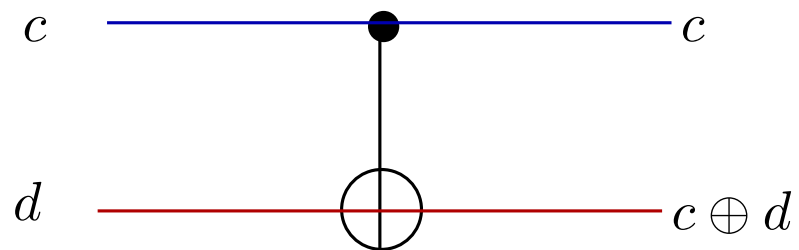
CNOT versus $X \otimes Y$

If we start with $\frac{1}{\sqrt{2}}(|00\rangle + |11\rangle)$, we get

- $(X \otimes Y) \left(\frac{1}{\sqrt{2}}(|00\rangle + |11\rangle) \right) = \frac{i}{\sqrt{2}}(-|00\rangle + |11\rangle)$.
- $CNOT\left(\frac{1}{\sqrt{2}}(|00\rangle + |11\rangle)\right) = \frac{1}{\sqrt{2}}(|00\rangle + |10\rangle)$

The input state was **entangled**. The entanglement remains after the transformation $X \otimes Y$, but it is not there after CNOT.

Graphical Notation: CNOT



CNOT

Controlled U Gate

For every 1-qubit transformation U we can have a **controlled- U gate**, U^c .

$$U^c |ab\rangle = \begin{cases} |ab\rangle & \text{if } a = 0, \\ |a(Ub)\rangle & \text{if } a = 1. \end{cases}$$

So $U^c = |0\rangle\langle 0| \otimes I + |1\rangle\langle 1| \otimes U$.

Controlled- U Transformation

- Let $U = \begin{bmatrix} u_{11} & u_{12} \\ u_{21} & u_{22} \end{bmatrix}$ be a 1-qubit unitary transformation.
- The transformation matrix for

$$U^c = \begin{bmatrix} 1 & 0 & 0 & 0 \\ 0 & 1 & 0 & 0 \\ 0 & 0 & u_{11} & u_{12} \\ 0 & 0 & u_{21} & u_{22} \end{bmatrix}.$$

Z^c, H and CNOT

We know that

$$H = \frac{1}{\sqrt{2}} \begin{bmatrix} 1 & 1 \\ 1 & -1 \end{bmatrix} = \frac{1}{\sqrt{2}} (|0\rangle\langle 0| + |0\rangle\langle 1| + |1\rangle\langle 0| - |1\rangle\langle 1|)$$

and $I = |0\rangle\langle 0| + |1\rangle\langle 1|$. So,

$$\begin{aligned} I \otimes H &= (|0\rangle\langle 0| + |1\rangle\langle 1|) \otimes \\ &\quad \frac{1}{\sqrt{2}} (|0\rangle\langle 0| + |0\rangle\langle 1| + |1\rangle\langle 0| - |1\rangle\langle 1|) \\ &= \frac{1}{\sqrt{2}} (|00\rangle\langle 00| + |00\rangle\langle 01| + |01\rangle\langle 00| - |01\rangle\langle 01| + \\ &\quad |10\rangle\langle 10| + |10\rangle\langle 11| + |11\rangle\langle 10| - |11\rangle\langle 11|) \end{aligned}$$

Z^c, H and CNOT

$$I \otimes H = \frac{1}{\sqrt{2}} \begin{bmatrix} 1 & 1 & 0 & 0 \\ 1 & -1 & 0 & 0 \\ 0 & 0 & 1 & 1 \\ 0 & 0 & 1 & -1 \end{bmatrix}, \quad Z^c = \begin{bmatrix} 1 & 0 & 0 & 0 \\ 0 & 1 & 0 & 0 \\ 0 & 0 & 1 & 0 \\ 0 & 0 & 0 & -1 \end{bmatrix}$$

So $(I \otimes H)Z^c(I \otimes H)$ is

$$\frac{1}{2} \begin{bmatrix} 1 & 1 & 0 & 0 \\ 1 & -1 & 0 & 0 \\ 0 & 0 & 1 & 1 \\ 0 & 0 & 1 & -1 \end{bmatrix} \begin{bmatrix} 1 & 0 & 0 & 0 \\ 0 & 1 & 0 & 0 \\ 0 & 0 & 1 & 0 \\ 0 & 0 & 0 & -1 \end{bmatrix} \begin{bmatrix} 1 & 1 & 0 & 0 \\ 1 & -1 & 0 & 0 \\ 0 & 0 & 1 & 1 \\ 0 & 0 & 1 & -1 \end{bmatrix}$$

Z^c, H and CNOT

So we have

$$\begin{aligned}
 (I \otimes H)Z^c(I \otimes H) &= \frac{1}{2} \begin{bmatrix} 1 & 1 & 0 & 0 \\ 1 & -1 & 0 & 0 \\ 0 & 0 & 1 & 1 \\ 0 & 0 & 1 & -1 \end{bmatrix} \begin{bmatrix} 1 & 1 & 0 & 0 \\ 1 & -1 & 0 & 0 \\ 0 & 0 & 1 & 1 \\ 0 & 0 & -1 & 1 \end{bmatrix} \\
 &= \frac{1}{2} \begin{bmatrix} 2 & 0 & 0 & 0 \\ 0 & 2 & 0 & 0 \\ 0 & 0 & 0 & 2 \\ 0 & 0 & 2 & 0 \end{bmatrix} = \begin{bmatrix} 1 & 0 & 0 & 0 \\ 0 & 1 & 0 & 0 \\ 0 & 0 & 0 & 1 \\ 0 & 0 & 1 & 0 \end{bmatrix} \\
 &= \text{CNOT}
 \end{aligned}$$

CNOT and Basis

- The 2-qubit CNOT gate behaves very similar to 2-bit Boolean gate, where the **control bit** remains unchanged and the other bit flips when the control bit is $|1\rangle$. This happens when the input state is in **standard basis**.
- But if the input state is not in standard basis, CNOT behaves differently.

CNOT On Hadamard Basis

The bases of a 2-qubit state space in Hadamard basis is $\{|++\rangle, |+-\rangle, |-+\rangle, |--\rangle\}$, where $|+\rangle = \frac{1}{\sqrt{2}}(|0\rangle + |1\rangle)$ and $|-\rangle = \frac{1}{\sqrt{2}}(|0\rangle - |1\rangle)$.

$$\begin{aligned} |++\rangle &= \frac{1}{\sqrt{2}}(|0\rangle + |1\rangle) \otimes \frac{1}{\sqrt{2}}(|0\rangle + |1\rangle) \\ &= \frac{1}{2} \begin{bmatrix} 1 \\ 1 \\ 1 \\ 1 \end{bmatrix} \otimes \begin{bmatrix} 1 \\ 1 \\ 1 \\ 1 \end{bmatrix} = \frac{1}{2} \begin{bmatrix} 1 \\ 1 \\ 1 \\ 1 \end{bmatrix}. \end{aligned}$$

CNOT On Hadamard Basis

Similarly,

$$|+-\rangle = \frac{1}{2} \begin{bmatrix} 1 \\ -1 \\ 1 \\ -1 \end{bmatrix}, | -+\rangle = \frac{1}{2} \begin{bmatrix} 1 \\ 1 \\ -1 \\ -1 \end{bmatrix}, |--\rangle = \frac{1}{2} \begin{bmatrix} 1 \\ -1 \\ -1 \\ 1 \end{bmatrix}.$$

CNOT On Hadamard Basis

$$CNOT |++\rangle = \begin{bmatrix} 1 & 0 & 0 & 0 \\ 0 & 1 & 0 & 0 \\ 0 & 0 & 0 & 1 \\ 0 & 0 & 1 & 0 \end{bmatrix} \begin{bmatrix} 1 \\ 1 \\ 1 \\ 1 \end{bmatrix} = \begin{bmatrix} 1 \\ 1 \\ 1 \\ 1 \end{bmatrix} = |++\rangle.$$

Similarly,

$$CNOT |+-\rangle = |--\rangle, CNOT |-+\rangle = |-+\rangle, CNOT |--\rangle = |+-\rangle.$$

The second qubit remains unchanged, the first qubit flips when the second one is $|-\rangle$.

Superdense Coding: An Application

Alice can transmit two classical bits of information to Bob by sending only one qubit.

- Initially, Alice has the first qubit and Bob has the second qubit of an entangled pair of qubits - $\frac{1}{\sqrt{2}}(|00\rangle + |11\rangle)$.
- Alice (Bob) can only transform her (his) qubit.

Superdense Coding: An Application

Alice encodes her classical bit pairs 00, 01, 10, 11 as follows and sends to Bob.

$$00 \mapsto (I \otimes I) \frac{1}{\sqrt{2}} (|00\rangle + |11\rangle) = \frac{1}{\sqrt{2}} (|00\rangle + |11\rangle),$$

$$01 \mapsto (X \otimes I) \frac{1}{\sqrt{2}} (|00\rangle + |11\rangle) = \frac{1}{\sqrt{2}} (|10\rangle + |01\rangle),$$

$$10 \mapsto (Z \otimes I) \frac{1}{\sqrt{2}} (|00\rangle + |11\rangle) = \frac{1}{\sqrt{2}} (|00\rangle - |11\rangle),$$

$$11 \mapsto (iY \otimes I) \frac{1}{\sqrt{2}} (|00\rangle + |11\rangle) = \frac{1}{\sqrt{2}} (-|10\rangle + |01\rangle).$$

Note that the second qubit is not touched. These transformations do not affect the entanglement.

Superdense Coding: An Application

After receiving the first qubit from Alice, Bob performs the following transformation on the entangled qubit pairs.

1. Applies **CNOT** that transformations the pair as follows:

$$\begin{aligned}\frac{1}{\sqrt{2}}(|00\rangle + |11\rangle) &\mapsto \frac{1}{\sqrt{2}}(|00\rangle + |10\rangle), \\ \frac{1}{\sqrt{2}}(|10\rangle + |01\rangle) &\mapsto \frac{1}{\sqrt{2}}(|11\rangle + |01\rangle), \\ \frac{1}{\sqrt{2}}(|00\rangle - |11\rangle) &\mapsto \frac{1}{\sqrt{2}}(|00\rangle + |10\rangle), \\ \frac{1}{\sqrt{2}}(-|10\rangle + |01\rangle) &\mapsto \frac{1}{\sqrt{2}}(-|11\rangle + |01\rangle).\end{aligned}$$

Superdense Coding: An Application

2. Applies $H \otimes I$:

$$\frac{1}{\sqrt{2}}(|00\rangle + |10\rangle) \mapsto |00\rangle,$$

$$\frac{1}{\sqrt{2}}(|11\rangle + |01\rangle) \mapsto |01\rangle,$$

$$\frac{1}{\sqrt{2}}(|00\rangle - |10\rangle) \mapsto |10\rangle,$$

$$\frac{1}{\sqrt{2}}(-|11\rangle + |01\rangle) \mapsto |11\rangle.$$

3. Bob measures the pair and recovers the two classical bits.

Superdense Coding: An Application

- The four qubit states produced by Alice are the **orthonormal Bell basis** - $\left\{ \frac{1}{\sqrt{2}}(|00\rangle + |11\rangle), \frac{1}{\sqrt{2}}(|10\rangle + |01\rangle), \frac{1}{\sqrt{2}}(|00\rangle - |11\rangle), \frac{1}{\sqrt{2}}(-|10\rangle + |01\rangle) \right\}$.
- So Bob can perform suitable measurement and identify them directly.

Superdense Coding: Note

- Two qubits are involved, but Alice does not use the other qubit.
- Any third party may supply the entangled qubits to Alice and Bob.

Teleportation: An Application

Teleportation in a sense is reverse of **superdense coding**.

Alice has a qubit in some unknown state $|x\rangle = a|0\rangle + b|1\rangle$, where $|a|^2 + |b|^2 = 1$. She wishes to transmit the state information to Bob using two Boolean bits through a classical channel, so that Bob can **reconstruct the qubit**.

Teleportation: An Application

- To start with, the First qubit of an entangled pair $|y\rangle = \frac{1}{\sqrt{2}}(|00\rangle + |11\rangle)$ is with Alice and the second qubit is with Bob.

- Alice starts with the 3-qubit state

$$\begin{aligned} |x\rangle \otimes |y\rangle &= (a|0\rangle + b|1\rangle) \otimes \frac{1}{\sqrt{2}}(|00\rangle + |11\rangle) \\ &= \frac{1}{\sqrt{2}}(a|000\rangle + a|011\rangle + b|100\rangle + b|111\rangle) \end{aligned}$$

- She can transform the first two qubits and Bob can transform the third qubit.

Teleportation: An Application

1. Alice applies $CNOT \otimes I$ on $|x\rangle \otimes |y\rangle$

$$\begin{aligned} & (CNOT \otimes I)(|x\rangle \otimes |y\rangle) \\ &= \frac{1}{\sqrt{2}}(CNOT \otimes I)(a|000\rangle + a|011\rangle + b|100\rangle + b|111\rangle) \\ &= \frac{1}{\sqrt{2}}(a|000\rangle + a|011\rangle + b|110\rangle + b|101\rangle). \end{aligned}$$

Teleportation: An Application

2. Then she applies $H \otimes I \otimes I$ on $(CNOT \otimes I)(|x\rangle \otimes |y\rangle)$ i.e.

$$\begin{aligned}
 & (H \otimes I \otimes I)(CNOT \otimes I)(|x\rangle \otimes |y\rangle) \\
 = & \frac{1}{\sqrt{2}}(H \otimes I \otimes I)(a|000\rangle + a|011\rangle + b|110\rangle + b|101\rangle) \\
 = & \frac{1}{2}(a(|000\rangle + |100\rangle + |011\rangle + |111\rangle) + \\
 & b(|010\rangle - |110\rangle + |001\rangle - |101\rangle)) \\
 = & \frac{1}{2}(|00\rangle(a|0\rangle + b|1\rangle) + |01\rangle(a|1\rangle + b|0\rangle) \\
 & + |10\rangle(a|0\rangle - b|1\rangle) + |11\rangle(a|1\rangle - b|0\rangle)).
 \end{aligned}$$

Teleportation: An Application

3. Alice measures the first two qubits in standard basis. The outcomes of measurement are $|00\rangle$, $|01\rangle$, $|10\rangle$, or $|11\rangle$ with equal probability.
4. Alice transmits two Boolean bits 00, 01, 10, 11 to Bob on classical channel, depending on the outcome of previous measurement.

Teleportation: An Application

As a result of Alice's measurement, the **projected state** of the third qubit of Bob is $(a|0\rangle + b|1\rangle)$, $(a|1\rangle + b|0\rangle)$, $(a|0\rangle - b|1\rangle)$, or $(a|1\rangle - b|0\rangle)$.

Teleportation: An Application

1. Bob receives the pair of bits and applies the following transformations on his qubit to bring it to the state of Alice's unknown qubit.

Boolean bits	Transformation
00	$I(a 0\rangle + b 1\rangle) = a 0\rangle + b 1\rangle$
01	$X(a 1\rangle + b 0\rangle) = a 0\rangle + b 1\rangle$
10	$Z(a 0\rangle - b 1\rangle) = a 0\rangle + b 1\rangle$
11	$iY(a 1\rangle - b 0\rangle) = a 0\rangle + b 1\rangle$

Controlled- U Transformation

- For every 1-qubit unitary transformation U , it is possible to implement a 2-qubit, **controlled- U** transformation, U^c , using **CNOT** gates and single qubit gates.
- We know that any single-qubit unitary transformation U can be decomposed as $e^{i\alpha}AXBXC$, where $ABC = I$.

Rotation gates

We know that

$$R_x(\alpha) = \begin{bmatrix} \cos\left(\frac{\alpha}{2}\right) & -i \sin\left(\frac{\alpha}{2}\right) \\ -i \sin\left(\frac{\alpha}{2}\right) & \cos\left(\frac{\alpha}{2}\right) \end{bmatrix}$$

$$R_y(\alpha) = \begin{bmatrix} \cos\left(\frac{\alpha}{2}\right) & -\sin\left(\frac{\alpha}{2}\right) \\ \sin\left(\frac{\alpha}{2}\right) & \cos\left(\frac{\alpha}{2}\right) \end{bmatrix}$$

$$R_z(\alpha) = \begin{bmatrix} \cos\left(\frac{\alpha}{2}\right) - i \sin\left(\frac{\alpha}{2}\right) & & 0 \\ & 0 & \cos\left(\frac{\alpha}{2}\right) + i \sin\left(\frac{\alpha}{2}\right) \\ & & & \end{bmatrix} \cdot$$

$$= \begin{bmatrix} e^{-i\frac{\alpha}{2}} & 0 \\ 0 & e^{i\frac{\alpha}{2}} \end{bmatrix} \cdot$$

Hadamard Gate an Example

Observe that

$$\begin{aligned}
 & e^{i\pi/2} R_z(0) R_y(\pi/2) R_z(\pi) \\
 = & i \begin{bmatrix} 1 & 0 \\ 0 & 1 \end{bmatrix} \begin{bmatrix} \cos \pi/4 & -\sin \pi/4 \\ \sin \pi/4 & \cos \pi/4 \end{bmatrix} \begin{bmatrix} e^{-i\pi/2} & 0 \\ 0 & e^{i\pi/2} \end{bmatrix} \\
 = & i \frac{1}{\sqrt{2}} \begin{bmatrix} 1 & -1 \\ 1 & 1 \end{bmatrix} \begin{bmatrix} -i & 0 \\ 0 & i \end{bmatrix} \\
 = & i \frac{1}{\sqrt{2}} \begin{bmatrix} -i & -i \\ -i & i \end{bmatrix} = H.
 \end{aligned}$$

Hadamard Gate an Example

So we have $H = e^{i\pi/2} A X B X C$, where

$$A = R_y(\pi/4),$$

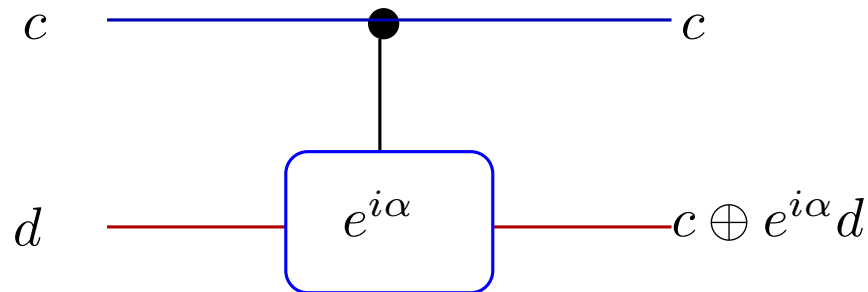
$$B = R_y(-\pi/4)R_z(-\pi/2),$$

$$C = R_z(\pi/2),$$

such that $ABC = I$.

Controlled-Phase Shift

If $U = e^{i\alpha} AXBXC$, where $ABC = I$, then in U^c the first operation is **controlled phase shift**, $(e^{i\alpha})^c$.



C-Phase-shift

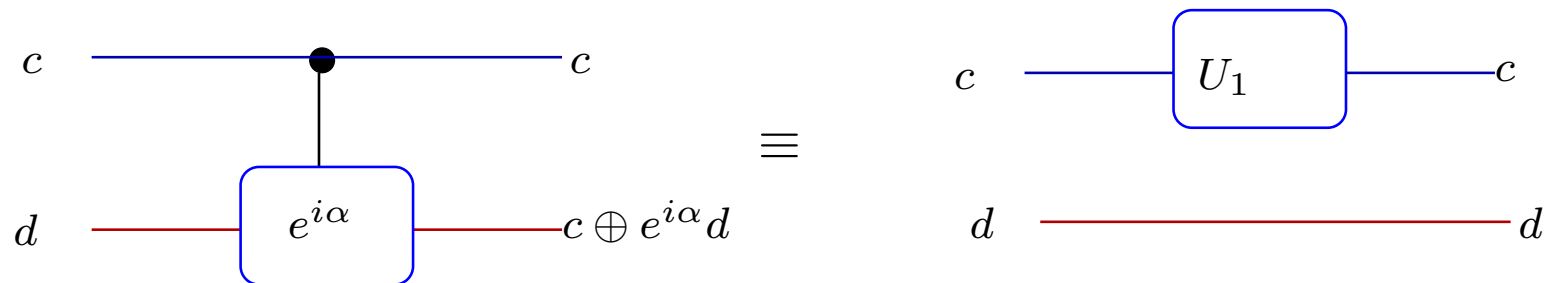
$|00\rangle \mapsto |00\rangle$, $|01\rangle \mapsto |01\rangle$, $|10\rangle \mapsto |1\rangle \otimes e^{i\alpha} |0\rangle$,
and $|11\rangle \mapsto |1\rangle \otimes e^{i\alpha} |1\rangle$.

Controlled-Phase Shift

- We observe that $|1\rangle \otimes e^{i\alpha} |x\rangle = e^{i\alpha} |1\rangle \otimes |x\rangle$, where $x \in \{0, 1\}$.
- We need a 1-qubit transform U_1 so that $U_1 |0x\rangle = |0x\rangle$ and $U_1 |1x\rangle = e^{i\alpha} |1\rangle \otimes |x\rangle$. So $(e^{i\alpha})^c$ is implemented as $U_1 \otimes I$, where

$$U_1 = \begin{bmatrix} 1 & 0 \\ 0 & e^{i\alpha} \end{bmatrix} = e^{i\alpha/2} I R_z(\alpha).$$

Controlled-Phase Shift

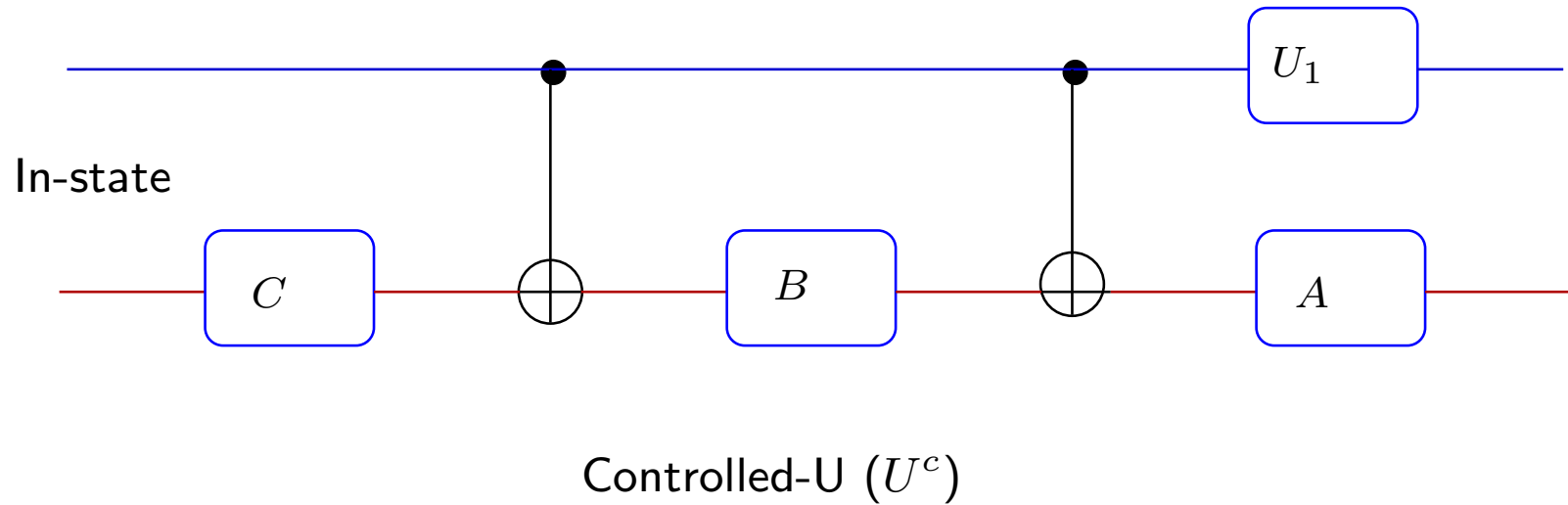


C-Phase-shift

Controlled- U Transformation

- If the control bit is $|1\rangle$, the state of the data bit is $Ud = (e^{i\alpha}AXBXC)d$.
- If the control bit is $|0\rangle$, the state of the data bit is $Id = (ABC)d$.
- The circuit is as follows:

Controlled- U Gate



Controlled- H Transformation: H^c

As we know

$$H = e^{i\pi/2} R_z(0) R_y(\pi/2) R_z(\pi) = e^{i\pi/2} AXBXC.$$

So we can construct **controlled-H** (H^c) gate.

Multi-qubit Control

- We can generalise the **single-control 2-qubit** unitary transformation to **multiply-controlled multi-qubit** unitary transformation.
- We have seen 3-bit reversible Boolean gates e.g. **Toffoli gate** and **Fredkin gate**, with two control-bits.

Multi-qubit Control

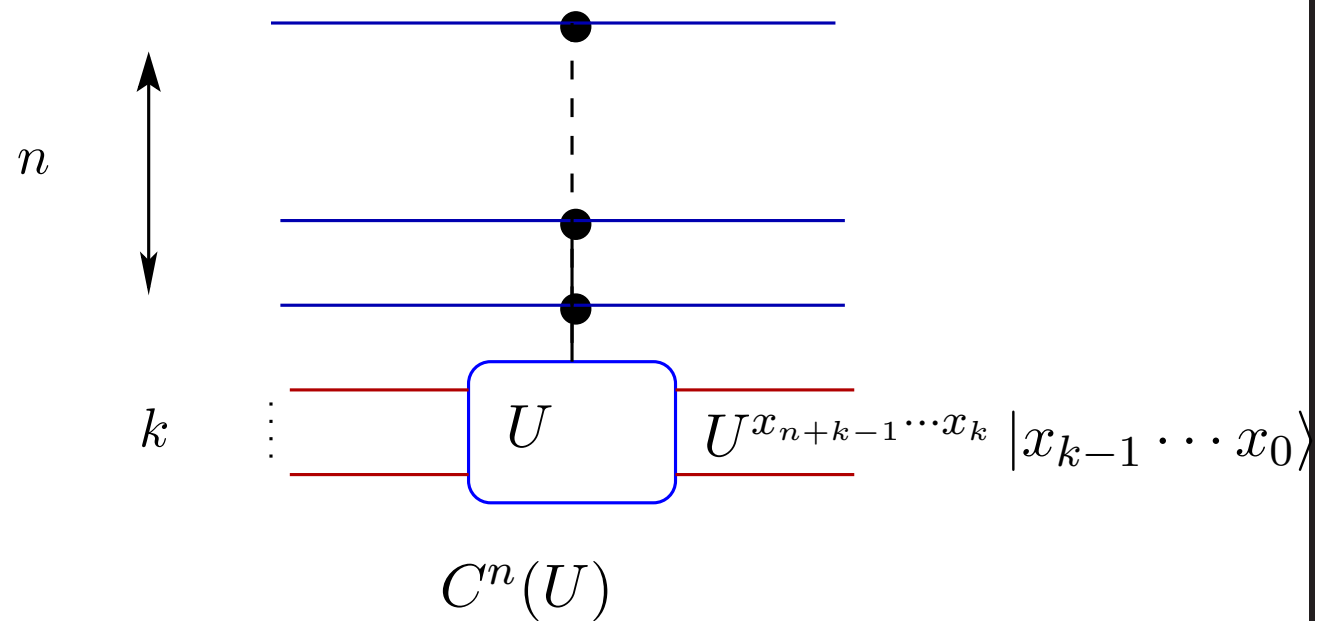
- Let U be a k -qubit unitary operator and there are n -control qubits.
- So we have a $(n + k)$ -qubits unitary operator $C^n(U)$ controlled by n -qubits.

$$\begin{aligned}
 & C^n(U) |x_{n+k-1} \cdots x_k\rangle |x_{k-1} \cdots x_0\rangle \\
 &= |x_{n+k-1} \cdots x_k\rangle U^{x_{n+k-1} \cdots x_k} |x_{k-1} \cdots x_0\rangle,
 \end{aligned}$$

U is applied on $|x_{k-1} \cdots x_0\rangle$ if

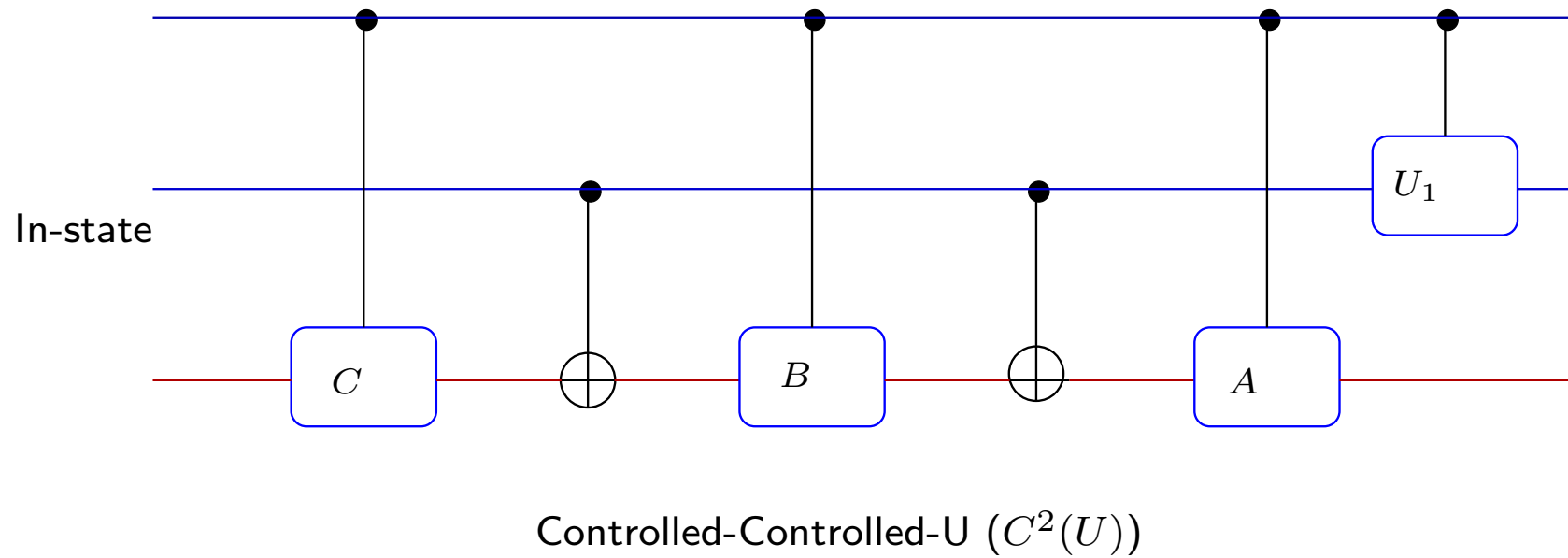
$$x_{n+k-1} = \cdots = x_k = 1.$$

Multi-Qubit Controlled Circuit



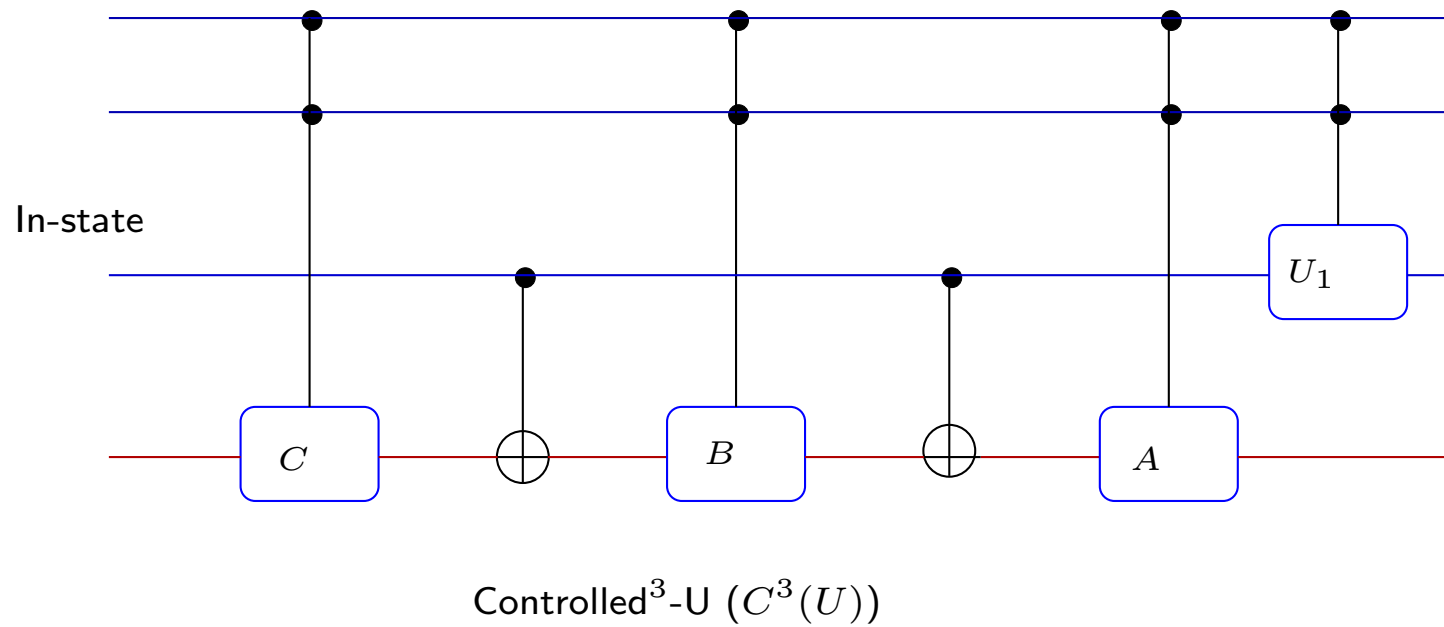
Multi-qubit Control

- We shall consider $k = 1$ and $n \geq 1$.
- The circuit for $n = 1$ can be used for $n = 2$ by replacing the 1-qubit gates A, B, C and U_1 by the corresponding control gates.

$C^2(U)$ Gate: Diagram

$C^2(U)$ Gate count

Each single-qubit control gate require **two CNOT** and **four single-qubit unitary** gates. So all together the requirement is $4^2 = 16$, single-qubit gates and $2 + 2 \cdot 4 = 10$, CNOT gates.

$C^3(U)$ Gate: Diagram

$C^k(U)$ Gate count

The number of 1-qubit gates are 4^k and the number of CNOT gates are

$$2 + 2 \cdot 4 + \dots + 2 \cdot 4^{k-1} = \frac{2}{3}(4^k - 1).$$

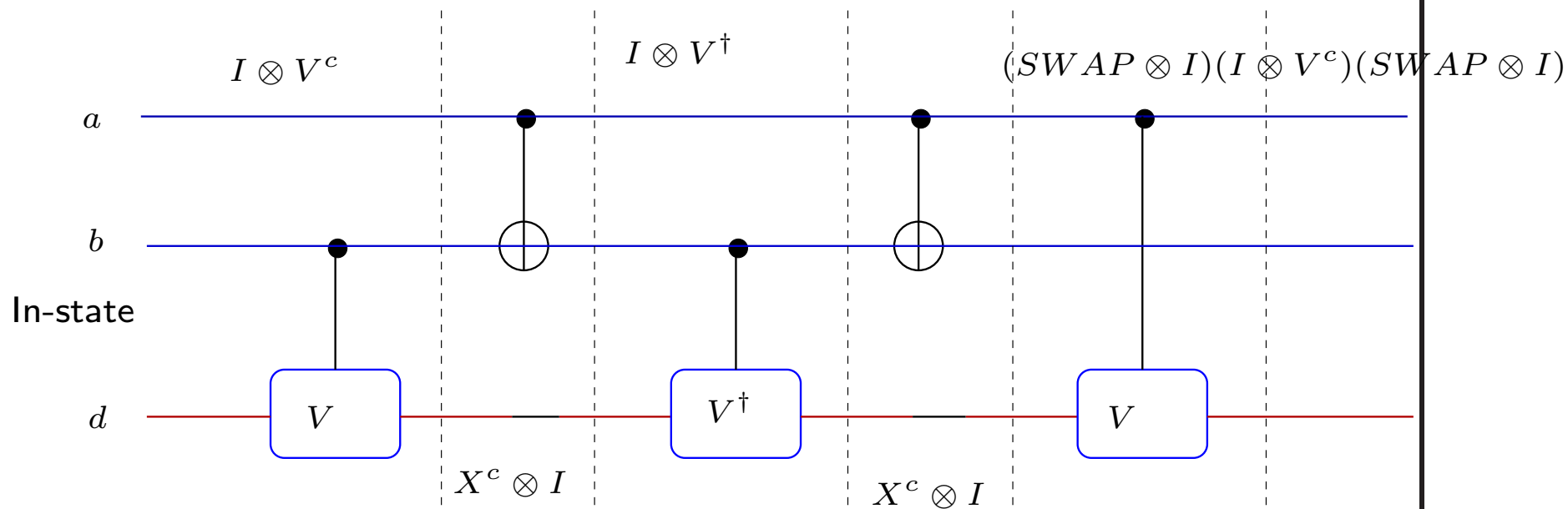
$$C^2(U) \text{ where } U = V^2$$

If the 1-qubit unitary operator $U = V^2$ where V is also unitary, then

$$C^2(U) = (SWAP \otimes I)(I \otimes V^c)(SWAP \otimes I)(X^c \otimes I) \\ (I \otimes (V^\dagger)^c)(X^c \otimes I)(I \otimes V^c)$$

This scheme uses $3 \times 4 = 12$ single-qubit gates and $3 \times 2 + 2 = 8$ CNOT gates.

$C^2(U)$ where $U = V^2$: Diagram



$C^2(U)$ where $U = V^2$

$$C^2(U) \text{ where } U = V^2$$

We apply the given sequence of transformations on $|00d\rangle$, $|01d\rangle$, $|10d\rangle$ and $|11d\rangle$, where $d \in \{0, 1\}$.

$$0. \quad |00d\rangle \xrightarrow{I \otimes V^c} |00d\rangle \xrightarrow{X^c \otimes I} |00d\rangle \xrightarrow{I \otimes (V^\dagger)^c} |00d\rangle \xrightarrow{X^c \otimes I} |00d\rangle \\ \xrightarrow{SWAP \otimes I} |00d\rangle \xrightarrow{I \otimes V^c} |00d\rangle \xrightarrow{SWAP \otimes I} |00d\rangle$$

$$1. \quad |01d\rangle \xrightarrow{I \otimes V^c} |01\rangle V |d\rangle \xrightarrow{X^c \otimes I} |01\rangle V |d\rangle \xrightarrow{I \otimes (V^\dagger)^c} |01\rangle V^\dagger V |d\rangle \\ = |01d\rangle \xrightarrow{X^c \otimes I} |01d\rangle \xrightarrow{SWAP \otimes I} |10d\rangle \xrightarrow{I \otimes V^c} |10d\rangle \xrightarrow{SWAP \otimes I} |01d\rangle$$

$$C^2(U) \text{ where } U = V^2$$

$$\begin{aligned}
 2. \quad & |10d\rangle \xrightarrow{I \otimes V^c} |10d\rangle \xrightarrow{X^c \otimes I} |11d\rangle \xrightarrow{I \otimes (V^\dagger)^c} |11\rangle V^\dagger |d\rangle \xrightarrow{X^c \otimes I} \\
 & |10\rangle V^\dagger |d\rangle \xrightarrow{SWAP \otimes I} |01\rangle V^\dagger |d\rangle \xrightarrow{I \otimes V^c} |01\rangle V V^\dagger |d\rangle = \\
 & |01d\rangle \xrightarrow{SWAP \otimes I} |10d\rangle
 \end{aligned}$$

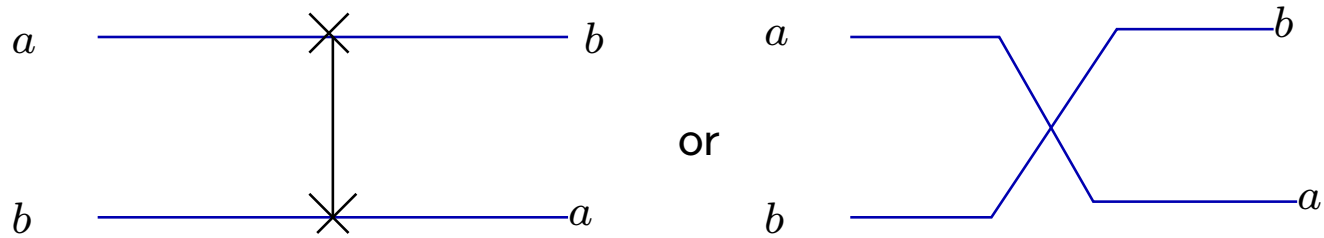
$$\begin{aligned}
 3. \quad & |11d\rangle \xrightarrow{I \otimes V^c} |11\rangle V |d\rangle \xrightarrow{X^c \otimes I} |10\rangle V |d\rangle \xrightarrow{I \otimes (V^\dagger)^c} |10\rangle V |d\rangle \\
 & \xrightarrow{X^c \otimes I} |11\rangle V |d\rangle \xrightarrow{SWAP \otimes I} |11\rangle V |d\rangle \xrightarrow{I \otimes V^c} |11\rangle V V |d\rangle = \\
 & |11\rangle U |d\rangle \xrightarrow{SWAP \otimes I} |11\rangle U |d\rangle
 \end{aligned}$$

SWAP Gate

An important 2-qubit gate is a **SWAP** gate. Its transition matrix is

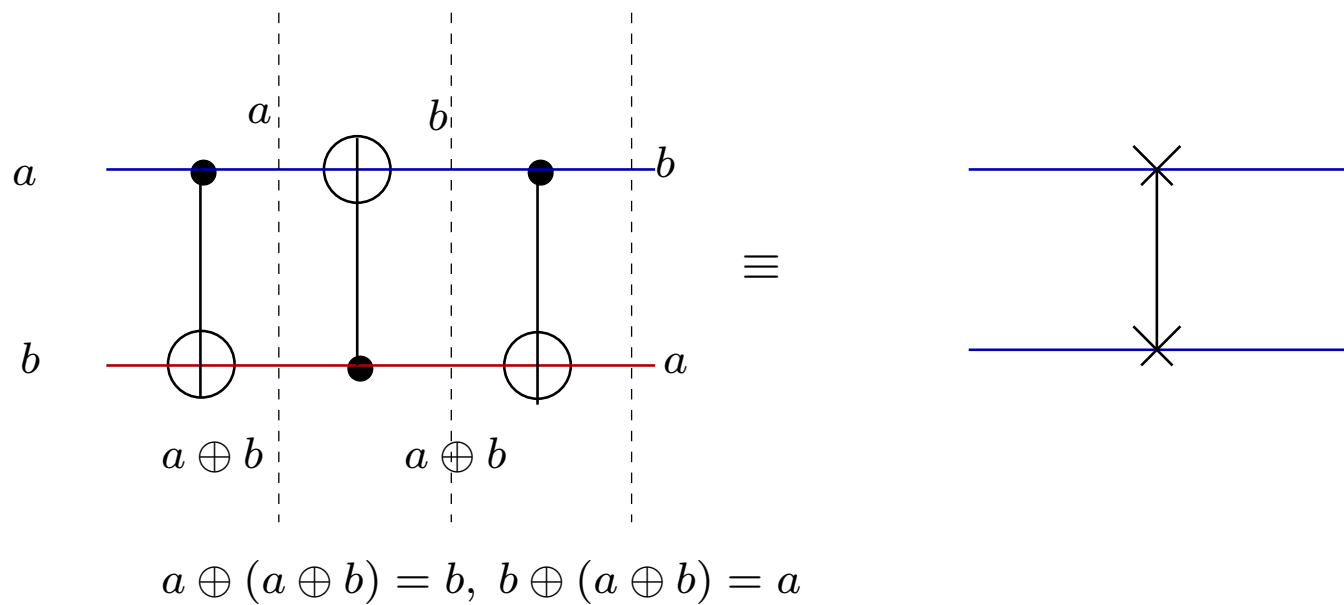
$$\begin{bmatrix} 1 & 0 & 0 & 0 \\ 0 & 0 & 1 & 0 \\ 0 & 1 & 0 & 0 \\ 0 & 0 & 0 & 1 \end{bmatrix}$$

SWAP Gate: Diagram



SWAP Gate

SWAP Gate using CNOT



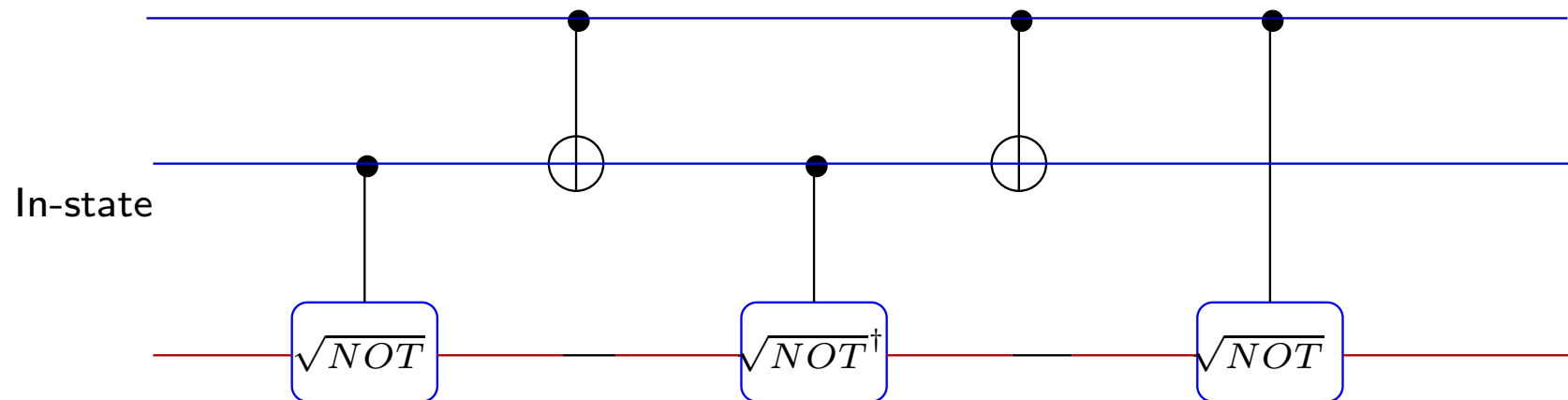
\sqrt{NOT} Gate

The following gate is known as \sqrt{NOT} such that $\sqrt{NOT} \cdot \sqrt{NOT} = NOT$.

$$\sqrt{NOT} = \frac{1}{2} \begin{bmatrix} 1+i & 1-i \\ 1-i & 1+i \end{bmatrix}.$$

This gate can be used to implement **CCNOT** or **Toffoli** gate.

Toffoli Gate using \sqrt{NOT} Gate



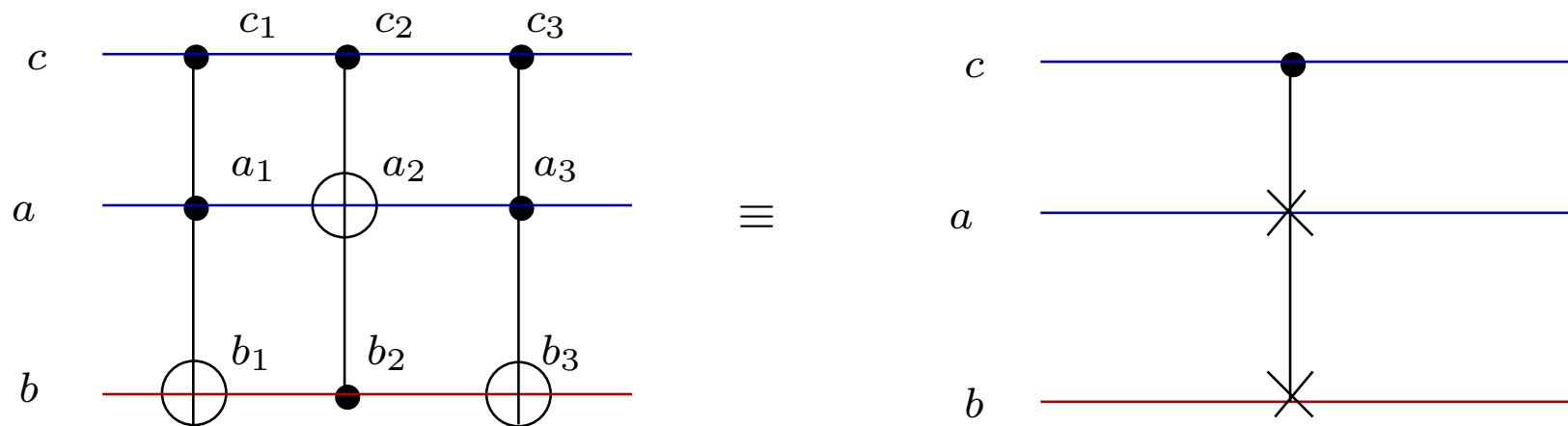
Note

Note that in Boolean logic, a **Toffoli gate** cannot be constructed using one-bit or two-bit gates. But a **quantum Toffoli** gate can be constructed using 2-qubit gates.

Fredkin or Controlled-Swap Gate

- We have already talked about the **Fredkin** or **controlled-SWAP** gate in connection to reversible Boolean logic.
- It is also known how a **SWAP** gate is implemented using 3 **CNOT** gates.
- So a Fredkin gate can be implemented as follows.

Fredkin or C-SWAP Gate using CCNOT



FREDKIN or Controlled-SWAP

Fredkin or C-SWAP Gate using CCCNOT

The computation of the left-hand circuit is

$$|c_1 a_1 b_1\rangle = |c, a, b \oplus ca\rangle$$

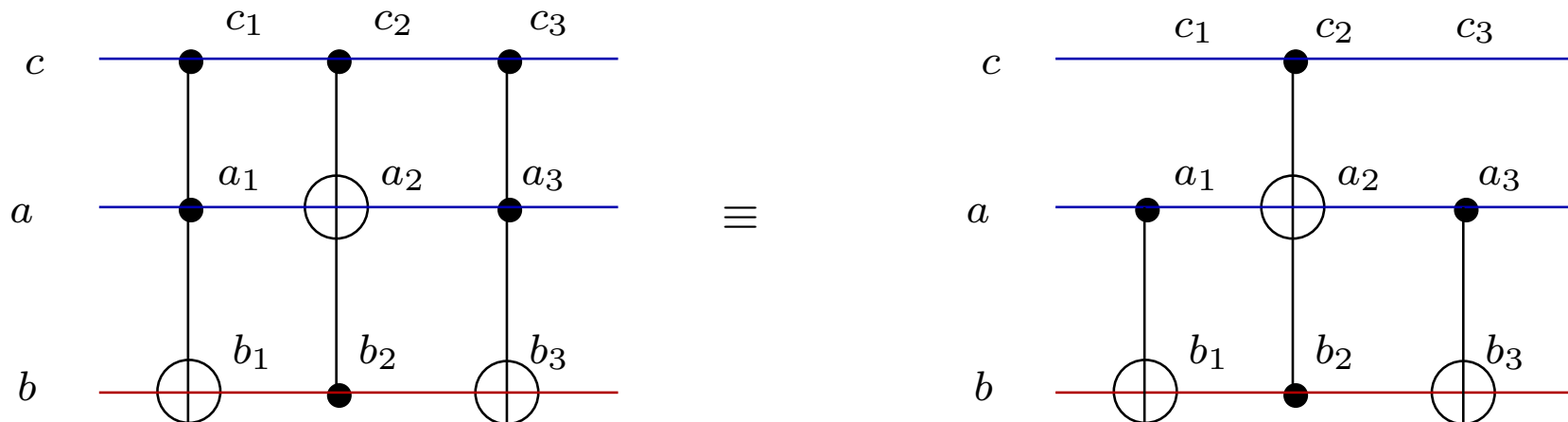
$$|c_2 a_2 b_2\rangle = |c, a \oplus c(b \oplus ca), b \oplus ca\rangle$$

$$|c_3 a_3 b_3\rangle = |c, a \oplus c(b \oplus ca), (b \oplus ca) \oplus c(a \oplus c(b \oplus ca))\rangle$$

- $c = 0$, $|c_3 a_3 b_3\rangle = |0, a, b\rangle$.
- $c = 1$, $|c_3 a_3 b_3\rangle = |1, b, a\rangle$ as $a \oplus (b \oplus a) = b$ and $(b \oplus a) \oplus (a \oplus (b \oplus a)) = a$.

Fredkin or Controlled-Swap Gate

We can replace the first and the third **CCNOT** gates by **CNOT** gates.



FREDKIN or Controlled-SWAP

The computation of the new circuit is

Fredkin Gate using CNOT and CCNOT

$$|c_1 a_1 b_1\rangle = |c, a, b \oplus a\rangle$$

$$|c_2 a_2 b_2\rangle = |c, a \oplus c(b \oplus a), b \oplus a\rangle$$

$$|c_3 a_3 b_3\rangle = |c, a \oplus c(b \oplus a), (b \oplus a) \oplus (a \oplus c(b \oplus a))\rangle .$$

- $c = 0$, $|c_3 a_3 b_3\rangle = |0, a, b\rangle$.
- $c = 1$, $|c_3 a_3 b_3\rangle = |1, b, a\rangle$ as $a \oplus (b \oplus a) = b$ and $(b \oplus a) \oplus (a \oplus (b \oplus a)) = a$.

Fredkin Gate using Only 2-Qubit Gates

- The **CCNOT** gate can be replaced by two **CNOT**, two \sqrt{NOT} and one \sqrt{NOT}^\dagger gate.
- So the Fredkin gate can be implemented using **seven 2-qubit** gates.
- This was impossible in classical Boolean logic.

H, S, T and CNOT

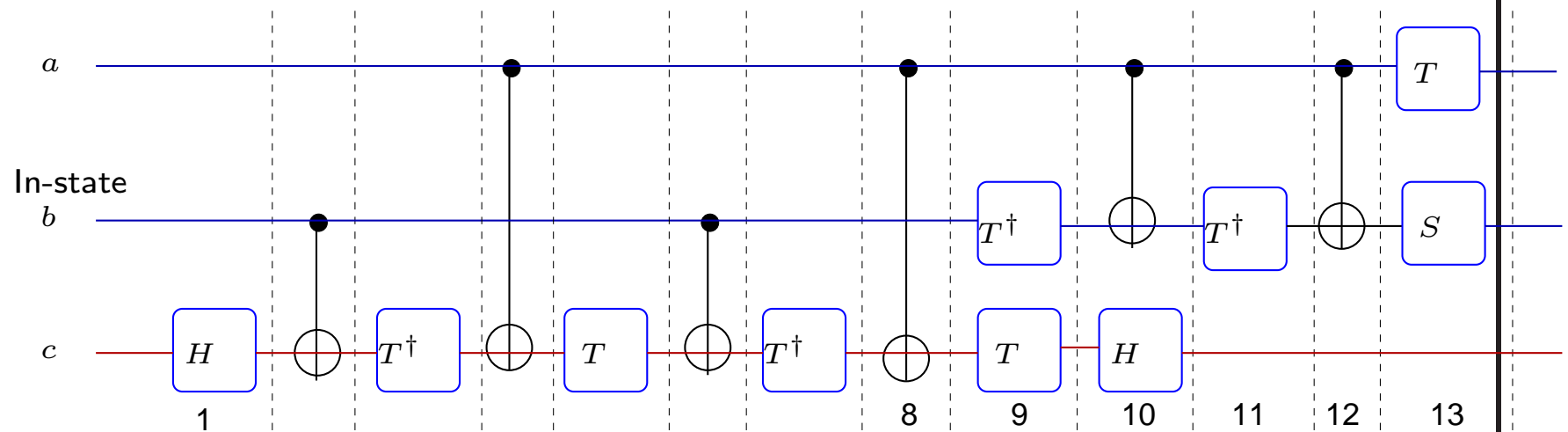
$$H = \frac{1}{\sqrt{2}} \begin{bmatrix} 1 & 1 \\ 1 & -1 \end{bmatrix}, S = \begin{bmatrix} 1 & 0 \\ 0 & i \end{bmatrix}, T = \begin{bmatrix} 1 & 0 \\ 0 & e^{i\pi/4} \end{bmatrix},$$

and

$$CNOT = \begin{bmatrix} 1 & 0 & 0 & 0 \\ 0 & 1 & 0 & 0 \\ 0 & 0 & 0 & 1 \\ 0 & 0 & 1 & 0 \end{bmatrix}$$

These four gates can be used to approximate any unitary transformation.

CCNOT using H , S , T and CNOT



CCNOT using H , S , T and CNOT

$$|abc\rangle \xrightarrow{1} |ab\rangle H |c\rangle \xrightarrow{2} \dots \xrightarrow{8} |ab\rangle X^a T^\dagger X^b T X^a T^\dagger X^b H |c\rangle ,$$

$$\xrightarrow{9} |a\rangle T^\dagger |b\rangle T X^a T^\dagger X^b T X^a T^\dagger X^b H |c\rangle$$

$$\xrightarrow{10} |a\rangle X^a T^\dagger |b\rangle H T X^a T^\dagger X^b T X^a T^\dagger X^b H |c\rangle$$

$$\xrightarrow{11} |a\rangle T^\dagger X^a T^\dagger |b\rangle H T X^a T^\dagger X^b T X^a T^\dagger X^b H |c\rangle$$

$$\xrightarrow{12} |a\rangle X^a T^\dagger X^a T^\dagger |b\rangle H T X^a T^\dagger X^b T X^a T^\dagger X^b H |c\rangle$$

$$\xrightarrow{13} T |a\rangle S X^a T^\dagger X^a T^\dagger |b\rangle H T X^a T^\dagger X^b T X^a T^\dagger X^b H |c\rangle$$

CCNOT using H , S , T and CNOT

- $a = 0$: $T|a\rangle = |0\rangle$, $SX^aT^\dagger X^aT^\dagger|b\rangle = S(T^\dagger)^2|b\rangle = |b\rangle$,
 $HTX^aT^\dagger X^bTX^aT^\dagger X^bH|c\rangle = |c\rangle$.
 So, $|0bc\rangle \xrightarrow{1\dots 12} |0bc\rangle$.
- $a = 1, b = 0$: $T|a\rangle = e^{i\pi/4}|1\rangle$, if we take the phase-factor $e^{i\pi/4}$ with the second term, we get
 $e^{i\pi/4}SX^aT^\dagger X^aT^\dagger|0\rangle = e^{i\pi/4}SXT^\dagger XT^\dagger|0\rangle = |0\rangle$.
 $HTX^aT^\dagger X^bTX^aT^\dagger X^bH|c\rangle = HTXT^\dagger TXT^\dagger H|c\rangle = |c\rangle$.
 So, $|10c\rangle \xrightarrow{1\dots 12} |10c\rangle$.

CCNOT using H , S , T and CNOT

- $a = 1 = b$: $T|a\rangle = e^{i\pi/4}|1\rangle$, if we take the phase-factor $e^{i\pi/4}$ with the second term, we get

$$e^{i\pi/4}SX^aT^\dagger X^aT^\dagger|1\rangle = e^{i\pi/4}SXT^\dagger XT^\dagger|1\rangle = i|1\rangle.$$

Transferring the phase-factor i to the third qubit state we get, $iHTX^aT^\dagger X^bTX^aT^\dagger X^bH|c\rangle =$

$$iH(TXT^\dagger X)(TXT^\dagger X)H|c\rangle = iH(-iZ)H|c\rangle =$$

$$HZH|c\rangle = X|c\rangle.$$

So, $|11c\rangle \xrightarrow{1\dots 12} |11\bar{c}\rangle.$
- So the circuit behaves like a **CCNOT** gate.

Controlled- U on $|0\rangle$

- The 1-qubit transformation U may be applied on the data-qubit when the control qubit is $|0\rangle$.
- The corresponding transformation matrix is

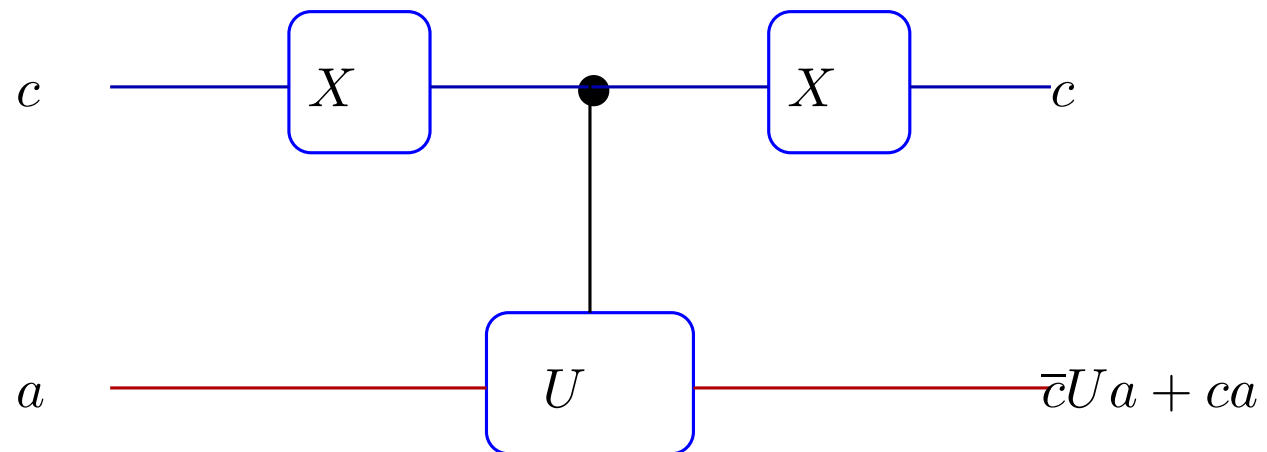
$$U_{|0\rangle}^c = \begin{bmatrix} u_{11} & u_{12} & 0 & 0 \\ u_{21} & u_{22} & 0 & 0 \\ 0 & 0 & 1 & 0 \\ 0 & 0 & 0 & 1 \end{bmatrix}$$

Controlled- U on $|0\rangle$

This can be achieved by $(X \otimes I) \circ U^c \circ (X \otimes I)$.

$$\begin{aligned}
 & \begin{bmatrix} 0 & 0 & 1 & 0 \\ 0 & 0 & 0 & 1 \\ 1 & 0 & 0 & 0 \\ 0 & 1 & 0 & 0 \end{bmatrix} \begin{bmatrix} 1 & 0 & 0 & 0 \\ 0 & 1 & 0 & 0 \\ 0 & 0 & u_{11} & u_{12} \\ 0 & 0 & u_{21} & u_{22} \end{bmatrix} \begin{bmatrix} 0 & 0 & 1 & 0 \\ 0 & 0 & 0 & 1 \\ 1 & 0 & 0 & 0 \\ 0 & 1 & 0 & 0 \end{bmatrix} \\
 = & \begin{bmatrix} u_{11} & u_{12} & 0 & 0 \\ u_{21} & u_{22} & 0 & 0 \\ 0 & 0 & 1 & 0 \\ 0 & 0 & 0 & 1 \end{bmatrix}
 \end{aligned}$$

Controlled- U on $|0\rangle$: Circuit



Controlled- U on $|0\rangle$