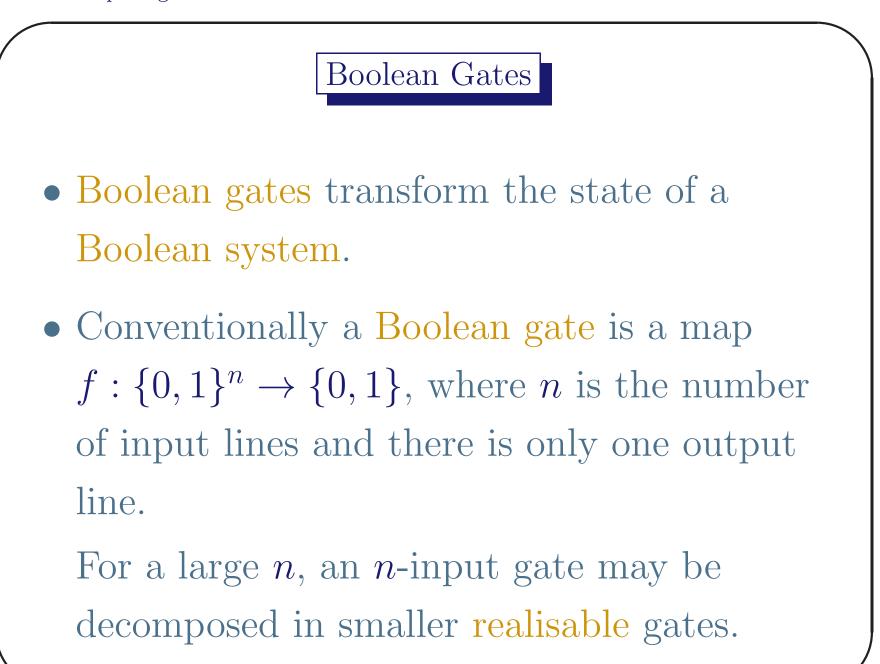
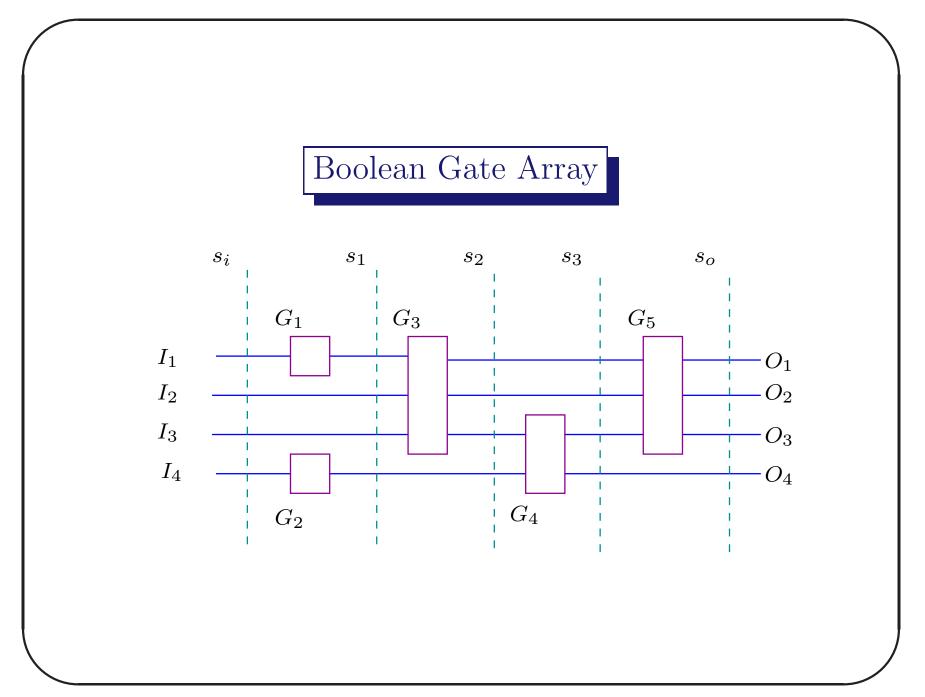
### **Boolean State Transformation**



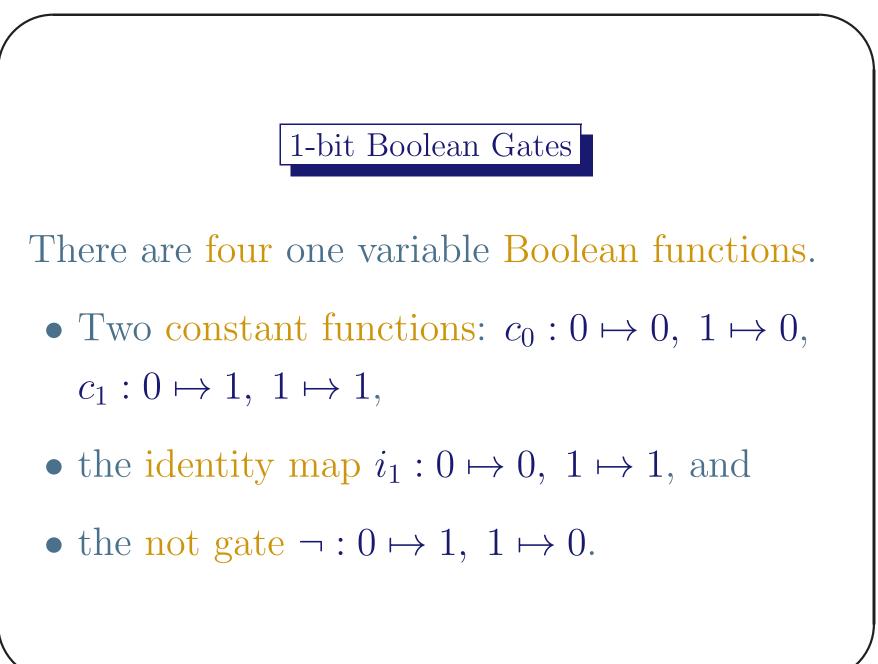


- But if the system state is represented by *n*-bits then a state transition map is g: {0,1}<sup>n</sup> → {0,1}<sup>n</sup>.
- The map g may be viewed as an n input and n output gate. This also may be realised using smaller gates.
- Following diagram is an example.





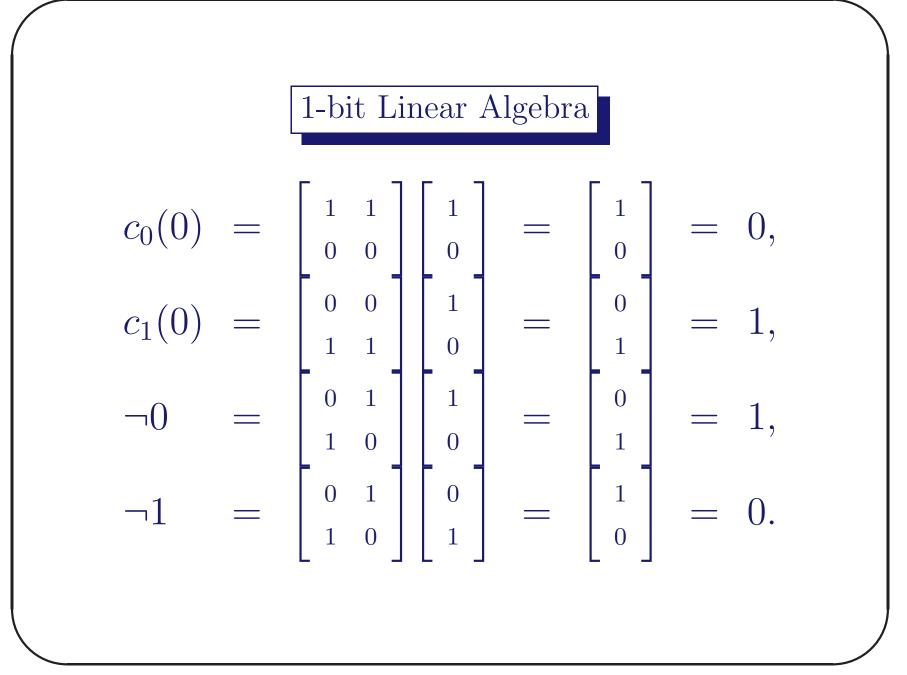
- $s_i$  is the input state,  $s_o$  is the output state.
- $s_1, s_2, s_3$  are intermediate states.
- Transition from s<sub>0</sub> to s<sub>1</sub> is through the gates
  G<sub>1</sub>, I, I, G<sub>2</sub>, where I may be viewed as
  identity map.
- This transition may be viewed as a 4-bit transformation  $G_1 \otimes I \otimes I \otimes G_2$ .
- Other transitions are similar.



1-bit Linear Algebra

If we encode Boolean 0 as  $\begin{bmatrix} 1\\0 \end{bmatrix}$  and 1 as  $\begin{bmatrix} 0\\1 \end{bmatrix}$ , the following transformation matrices represent the four gates.

$$c_{0} = \begin{bmatrix} 1 & 1 \\ 0 & 0 \end{bmatrix}, c_{1} = \begin{bmatrix} 0 & 0 \\ 1 & 1 \end{bmatrix}, i_{1} = \begin{bmatrix} 1 & 0 \\ 0 & 1 \end{bmatrix}, \neg : \begin{bmatrix} 0 & 1 \\ 1 & 0 \end{bmatrix}$$
  
The  $i_{1}$  and  $\neg$  gate are invertible, but other two are not.





- A 2 × 2 matrix over F<sub>2</sub> is a valid transformation matrix for a single-bit if every column has exactly one 1.
- This restriction is due to our encoding of 0 and 1.
- We get 2 × 2 = 4 valid transformation matrices corresponding to c<sub>0</sub>, c<sub>1</sub>, i<sub>1</sub> and ¬.
- Only two of them are reversible.

## Reversibility

- Reversibility of computation or invertibility of an operator is an issue even in classical computation.
- It is known from thermodynamics that there is no increase in entropy in a reversible process.
- But completely isentropic circuits are impossible to design.

## Reversibility

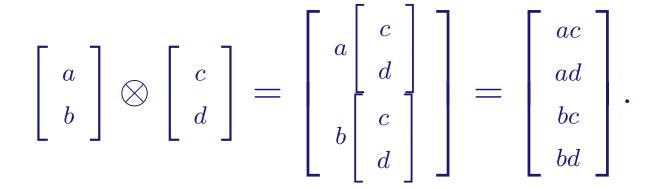
- There are adiabatic circuits that use reversible logic to reduce power consumption during switching.
- Landauer's principle claims that any "logically irreversible" change in information causes more change in entropy.
- A physical reversibility demands logical reversibility.

# Reversibility

- We shall see that state transition in a quantum mechanical system is reversible.
- Only the reversible classical logical gates may have quantum mechanical counterpart.
- So we explore the classical Boolean gates that are reversible as well as universal.



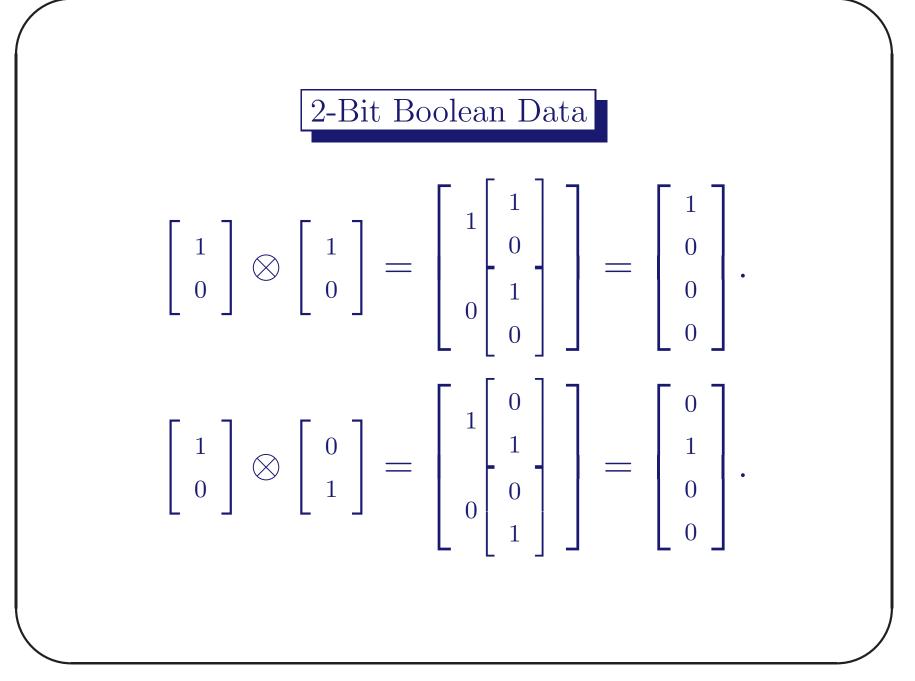
We define the tensor product of two 2-dimensional vectors  $\overrightarrow{x} = \begin{bmatrix} a \\ b \end{bmatrix}$  and  $\overrightarrow{y} = \begin{bmatrix} c \\ d \end{bmatrix}$ over some field F as

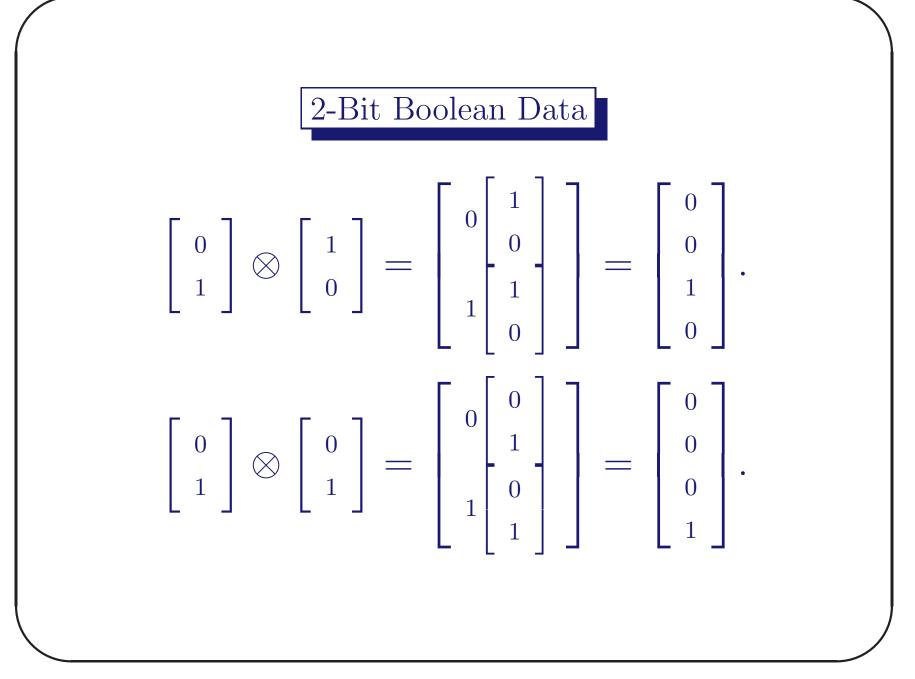


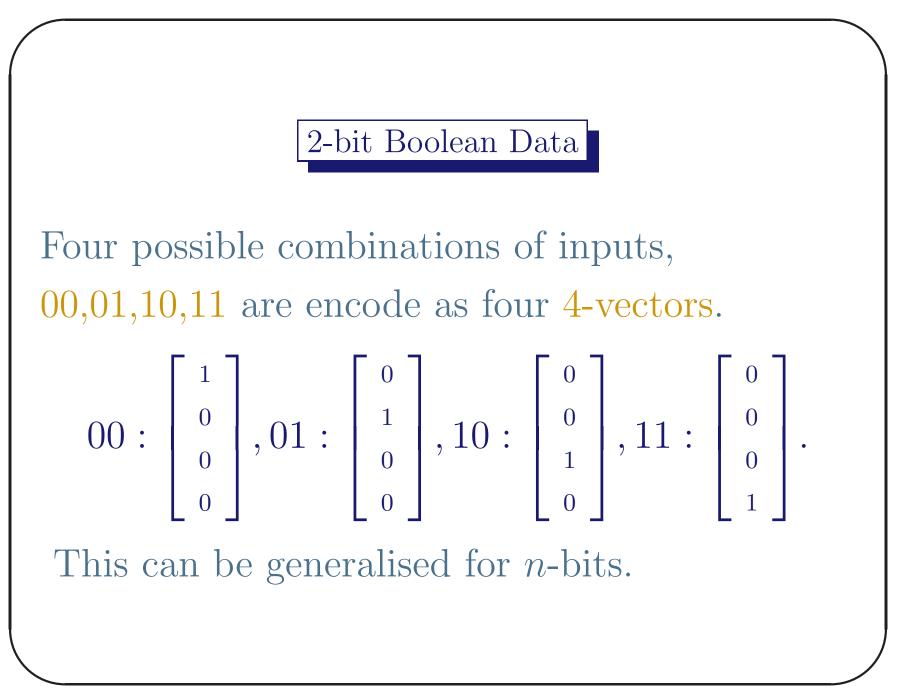
This can be generalised to higher dimensions.

### 2-Bit Boolean Data

1-bit Boolean data is encoded as  $0 = \begin{bmatrix} 1 \\ 0 \end{bmatrix}$  and  $1 = \begin{bmatrix} 0 \\ 1 \end{bmatrix}$ . Two bit Boolean data may be viewed as the tensor product of two 1-bit vectors.



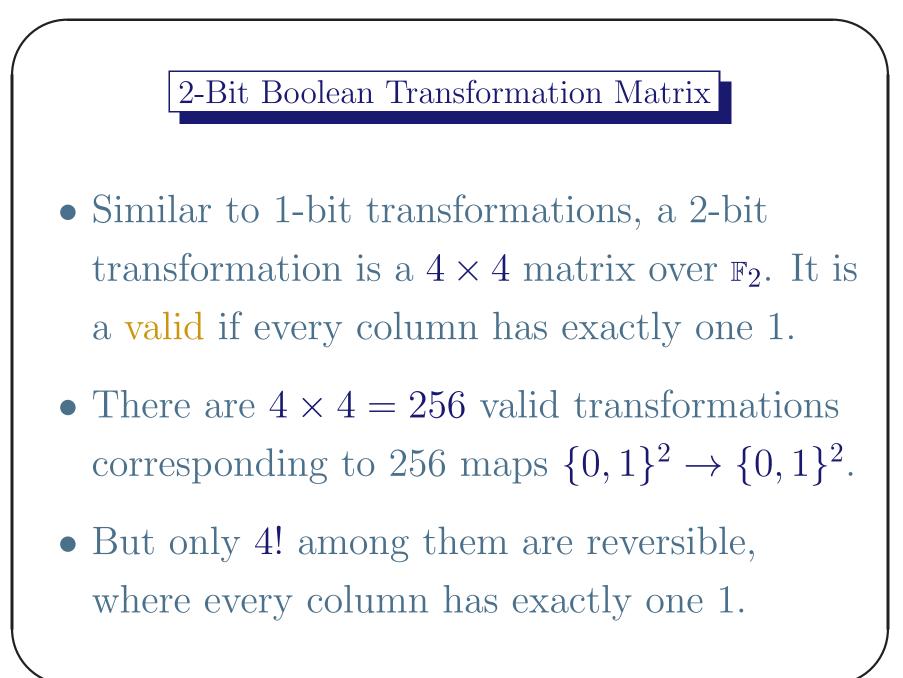


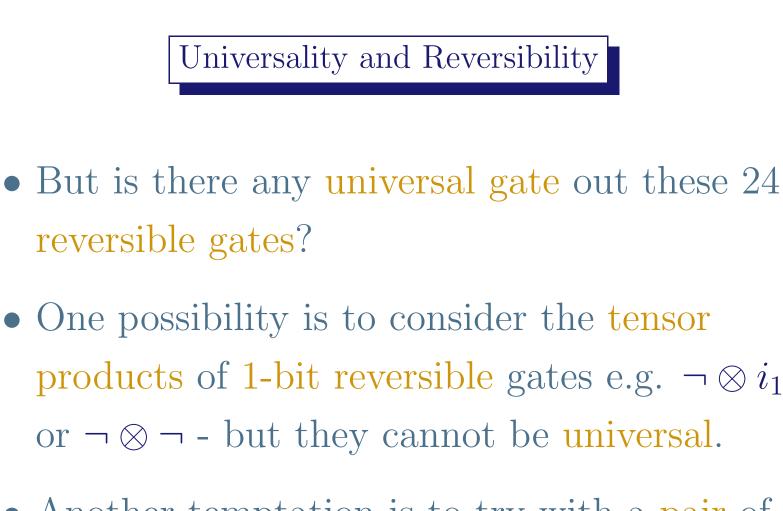


#### 2-bit Boolean Gates

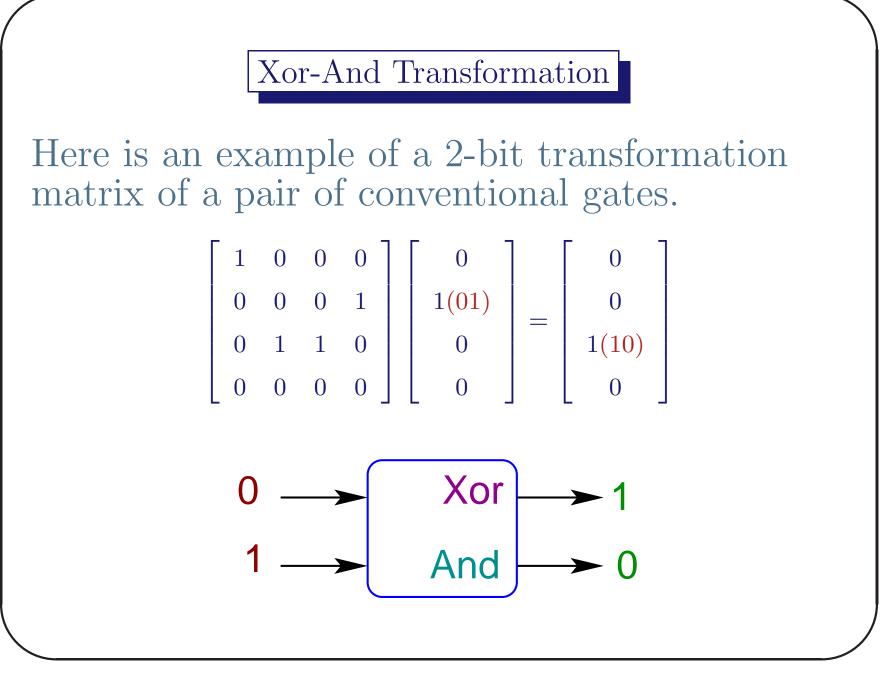
- There are sixteen 2-variable conventional Boolean functions e.g. and, nand, or, nor etc.
- They correspond to maps from  $\{0,1\}^2 \rightarrow \{0,1\}.$
- But we are interested about state transition maps from  $\{0,1\}^2 \rightarrow \{0,1\}^2$  that are reversible.

• There are total  $4^4 = 256$  such maps.





• Another temptation is to try with a pair of conventional gates and see what happens.



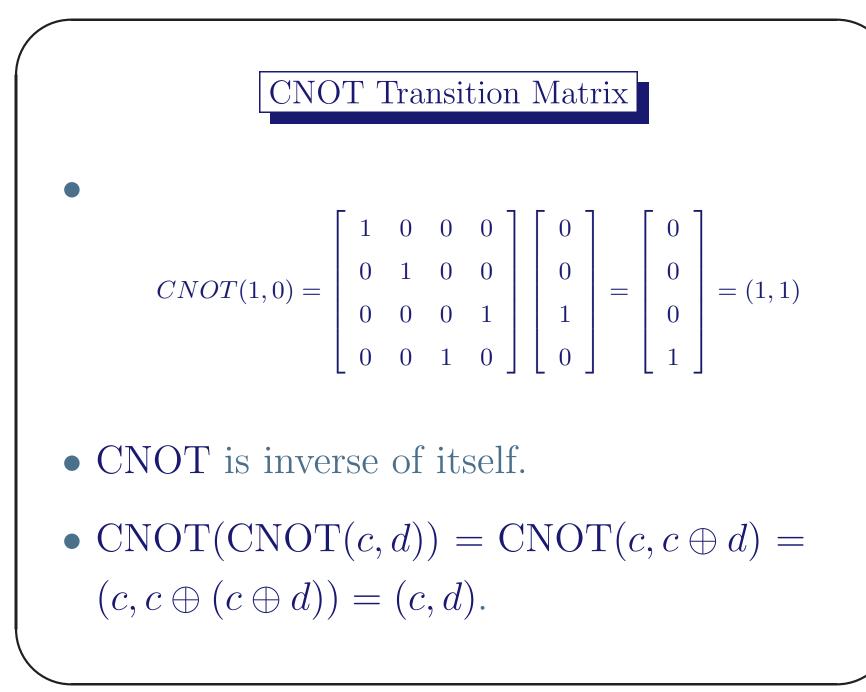


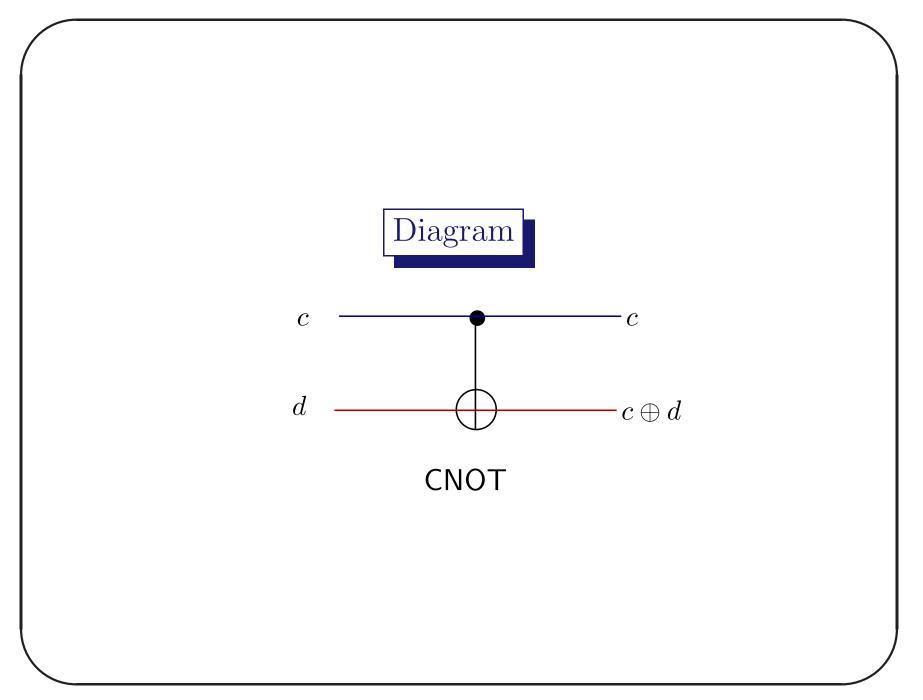
- The transformation matrix of Xor-And pair is not invertible - the  $4^{th}$ -row has all zeros.
- An invertible transformation matrix correspond to a permutation of rows of 4 × 4 identity matrix.

## CNOT Gate

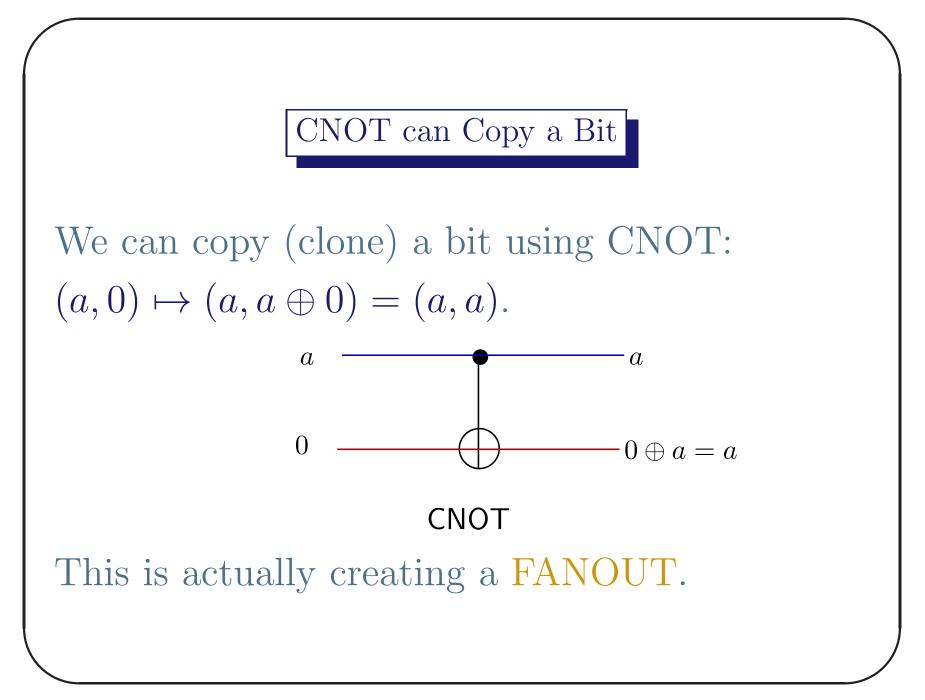
- A controlled not (CNOT) gate, has invertible transition matrix.
- Its inputs are (c, d), where c is the control input and d is the data input.
- The mapping is  $(c, d) \mapsto (c, c \oplus d)$ .

$$CNOT(c, d) = \begin{cases} (0, d) & \text{if } c = 0, \\ (1, \neg d) & \text{if } c = 1. \end{cases}$$



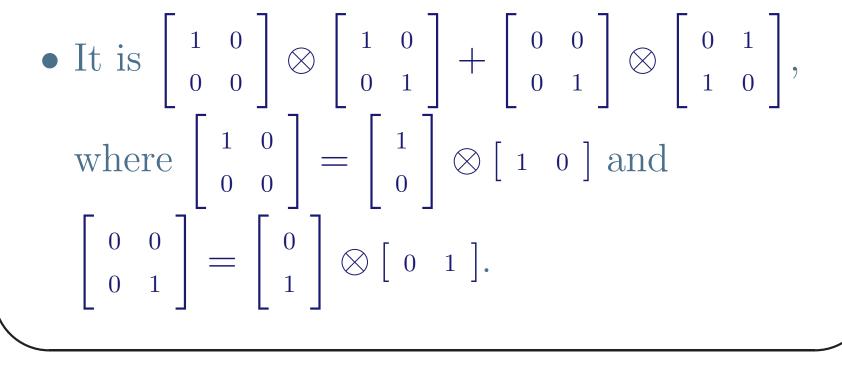


Lect 3





- Tensor products of 1-bit reversible gate are  $i_1 \otimes i_1, i_1 \otimes \neg, \neg \otimes i_1, \neg \otimes \neg$ .
- CNOT cannot be realised using them.



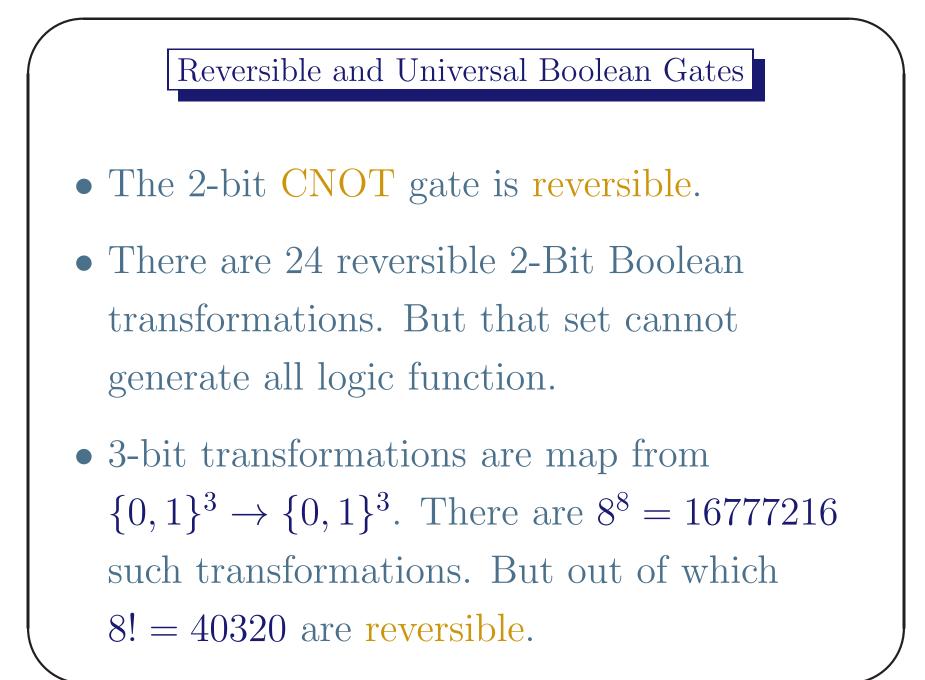
### 2-bit Universal Gate is Impossible

The truth-table for a 2-bit reversible gate is as follows:

$i_1$	$i_0$	<i>0</i> <sub>1</sub>	<i>O</i> <sub>0</sub>
0	0	$x_0$	$y_0$
0	1	$x_1$	$y_1$
1	0	$x_2$	$y_2$
1	1	$x_3$	$y_3$



- The output of a reversible gate is a permutation of {00,01,10,11}.
- It has exactly two 0's and two 1's per output column of the truth table.
- But i<sub>1</sub> ∧ i<sub>0</sub> has three 0's and i<sub>1</sub> ∨ i<sub>0</sub> has three
  1's in their output.
- None of them can be produced by a two bit reversible gate.



#### Universal Reversible Gates

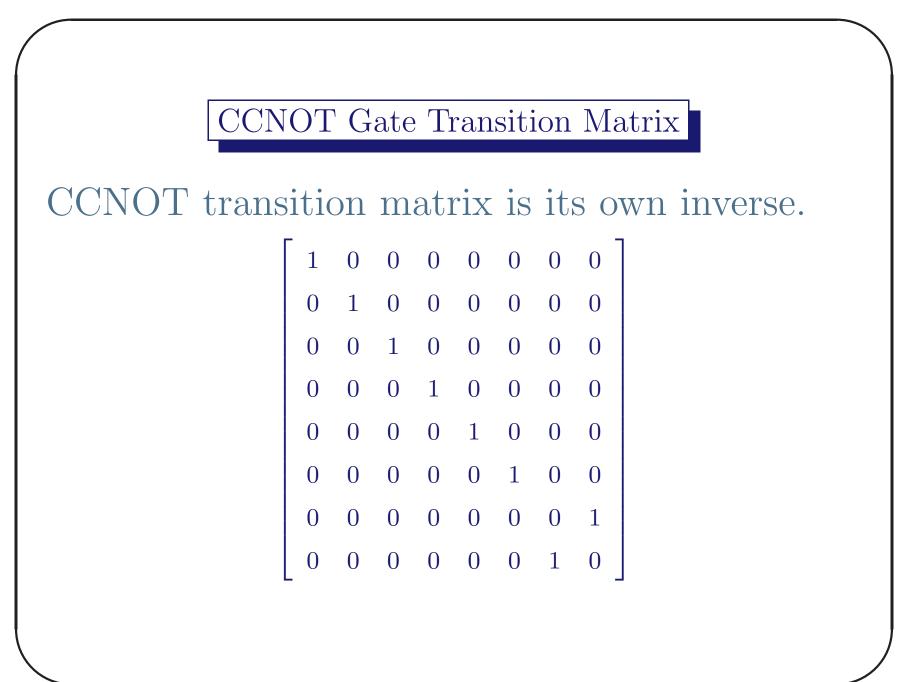
- One 3-bit reversible transformation is Toffoli gate or CCNOT gate. It was invented by Tommaso Toffoli from Italy.
- The Toffoli gate or CCNOT gate is known to be universal.

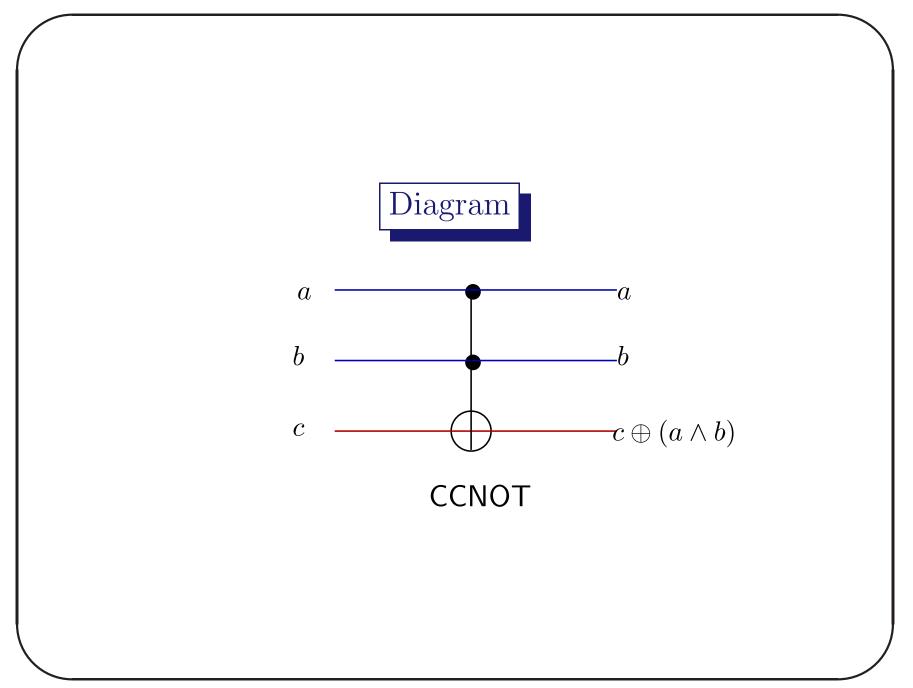


Following is the map of CCNOT gate:  $(x, y, c) \mapsto (x, y, c \oplus (x \land y))$ , where x, y remains unchanged,

$$c = \begin{cases} \neg c & \text{if } x = 1 = y, \\ c & \text{otherwise.} \end{cases}$$

If c = 0, the  $3^{rd}$  output is  $0 \oplus (x \land y) = x \land y$ .



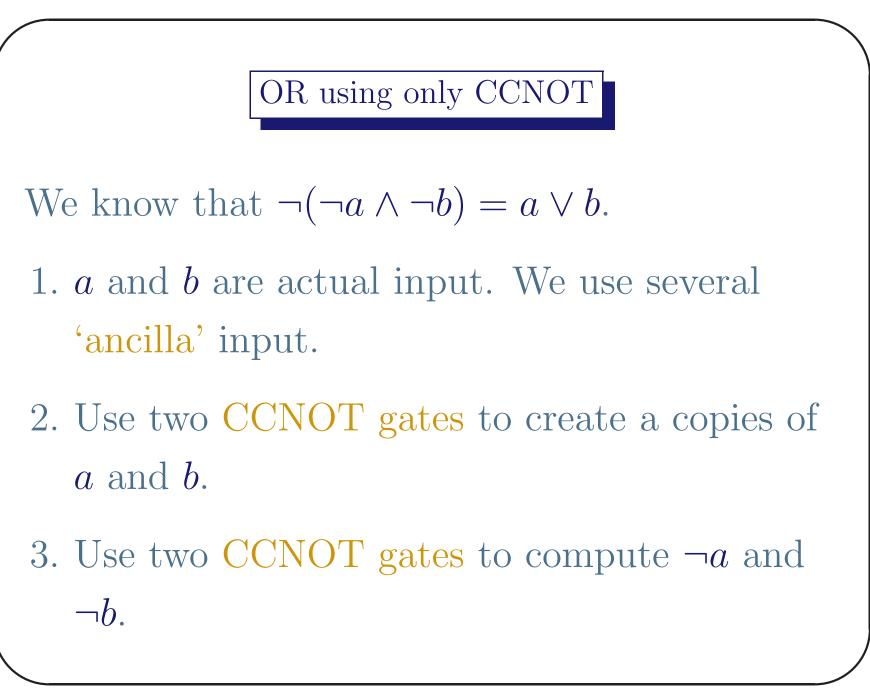




CCNOT gate can implement NOT, NAND and COPY/FANOUT operations if logic '0' and '1' are available.

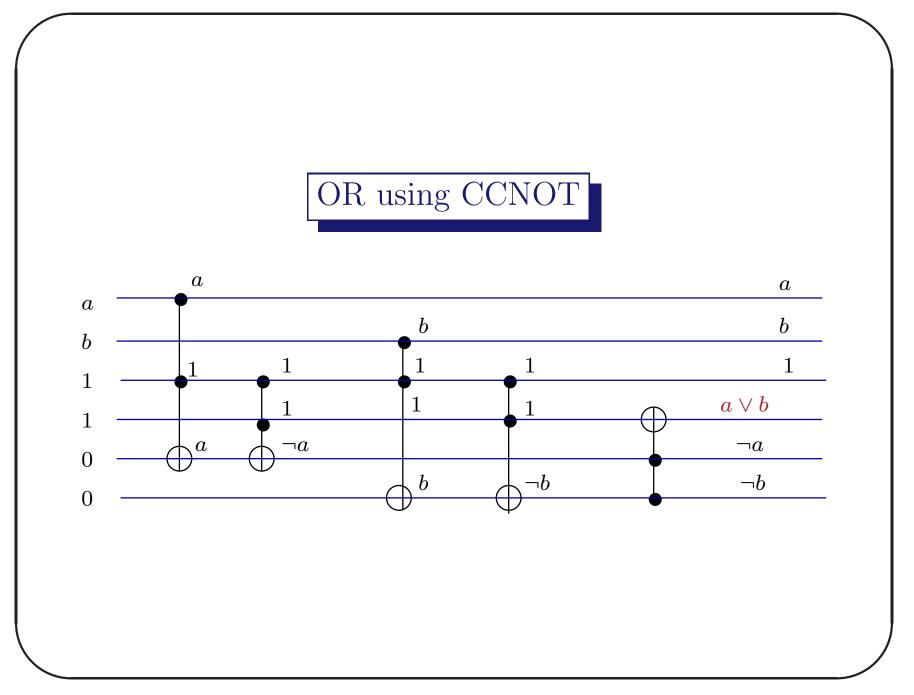
 $(1,1,c)\mapsto(1,1,\neg c),$ 

 $(a, b, 1) \mapsto (a, b, 1 \oplus (a \land b)) = (a, b, \neg (a \land b)),$  $(1, a, 0) \mapsto (1, a, 0 \oplus (1 \land a)) = (1, a, a).$ 





- 4. Finally another CCNOT gate to compute  $\neg(\neg a \land \neg b)$ .
- 5. We are not taking care of crossovers.



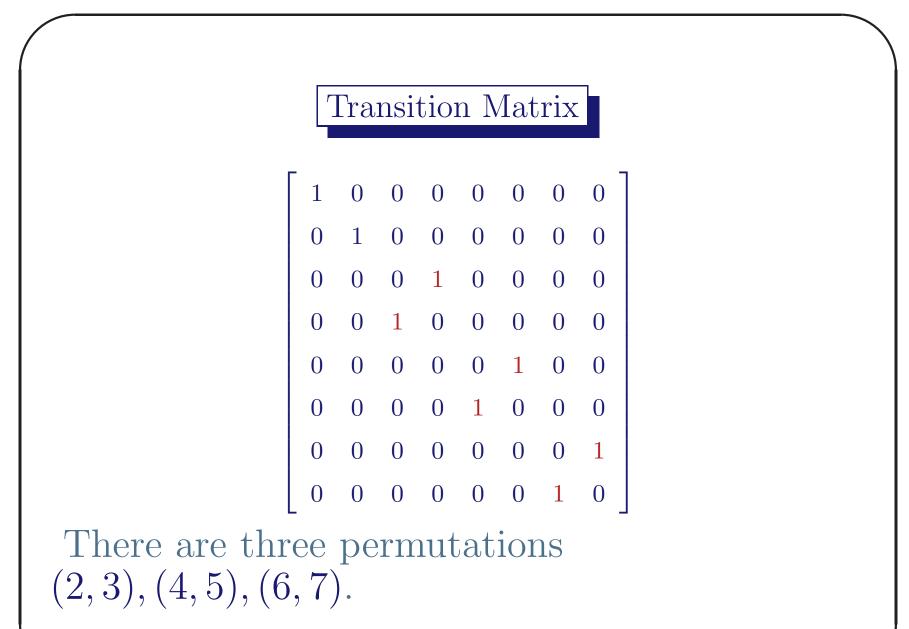


Note that we have not taken care of cross-over of bits in the diagram. There are several 'ancilla' inputs in state '0' and '1'. Also there are several useless outputs.



It is not very difficult to create a 3-bit reversible transformation that will compute a 2-bit function e.g. or when the  $3^{rd}$ -bit has a fixed value. Look at the following truth table.

	Truth Table							
	a	b	С	<i>o</i> <sub>1</sub>	<i>0</i> 2	<i>O</i> 3		
	0	0	0	0	0	0		
	0	0	1	0	0	1		
	0	1	0	0	1	1		
	0	1	1	0	1	0		
	1	0	0	1	0	1		
	1	0	1	1	0	0		
	1	1	0	1	1	1		
	1	1	1	1	1	0		
When $c = 0$ , $o_3 = a \lor b$ . Other bits are filled to maintain reversibility.								

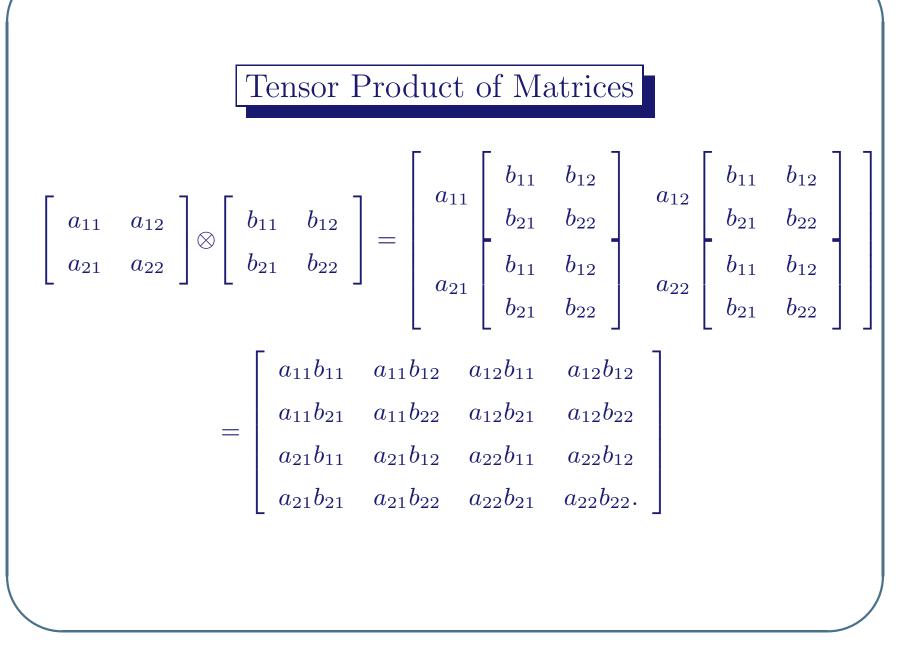


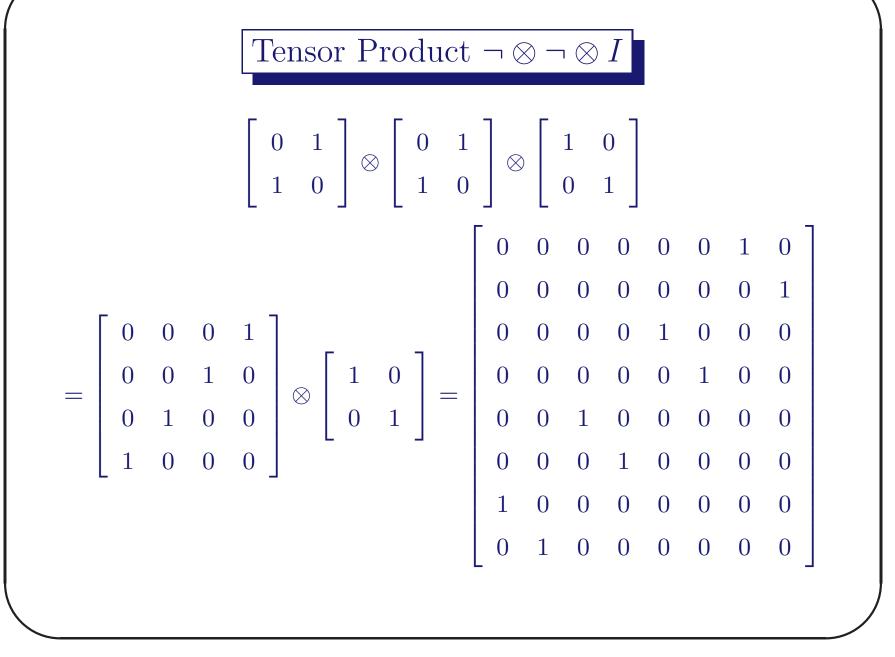


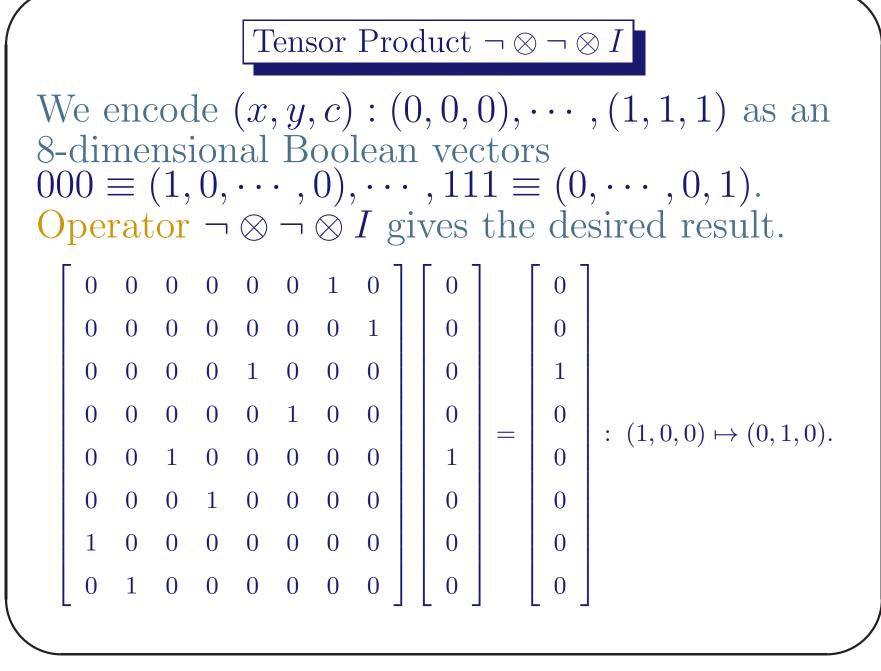
- Another way of getting OR using CCNOT and NOT are as follows.
- We negate the input a, b and keep c unchanged.
- This can be done by applying the transformation

$$\left[\begin{array}{cc} 0 & 1 \\ 1 & 0 \end{array}\right] \otimes \left[\begin{array}{cc} 0 & 1 \\ 1 & 0 \end{array}\right] \otimes \left[\begin{array}{cc} 1 & 0 \\ 0 & 1 \end{array}\right]$$

to a 3-bit vector.

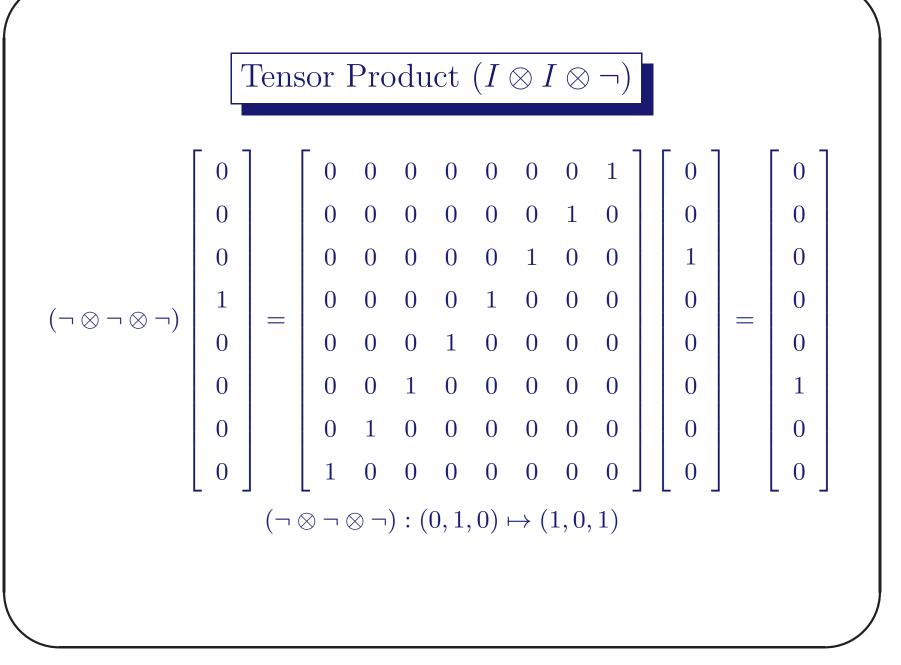






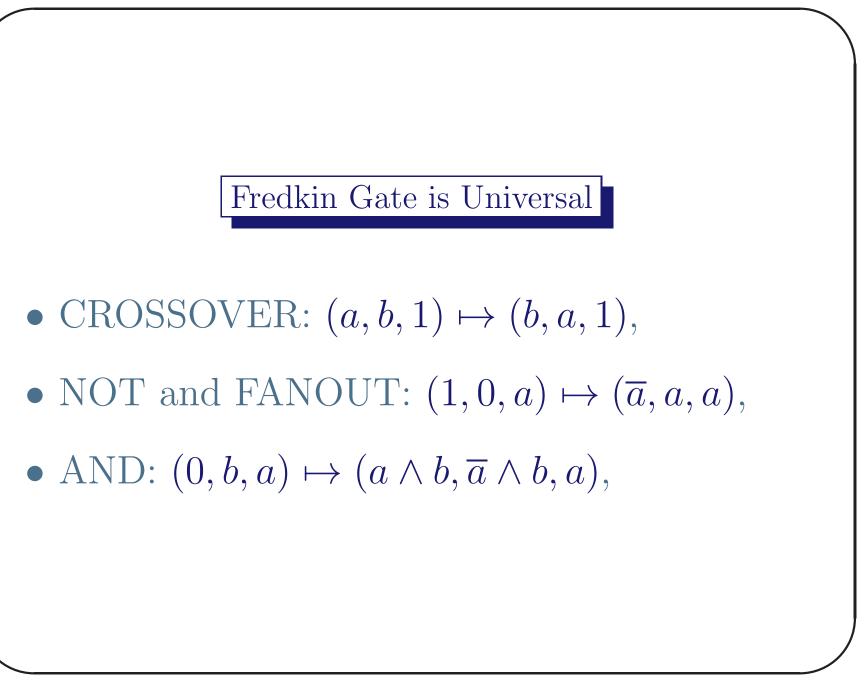


- The transformation  $(\neg \otimes \neg \otimes I)$  is clearly reversible (universal property).
- $CCNOT \circ (\neg \otimes \neg \otimes I)$  transforms  $(x, y, c) \mapsto (\neg x, \neg y, c \oplus (\neg x \land \neg y)).$
- We can apply  $\neg \otimes \neg \otimes \neg$  on the result. This gives us  $x \lor y$  when c = 0.

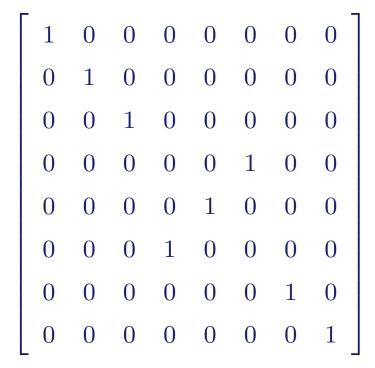


## Fredkin Gate

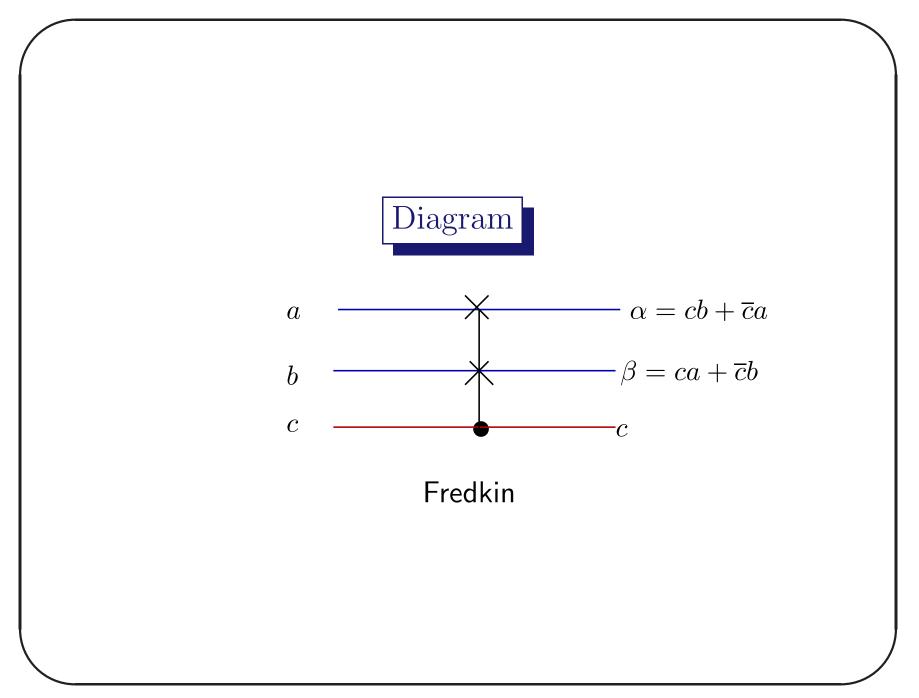
- Another important 3-input reversible gate is the Fredkin gate.
- The input-output relation is  $(x, y, c) \mapsto (cy + \overline{c}x, cx + \overline{c}y, c)$  i.e.  $(x, y, 0) \mapsto (x, y, 0)$  and  $(x, y, 1) \mapsto (y, x, 1)$ .
- If c = 0, x, y remains unchanged. If c = 1, the outputs are interchanged.





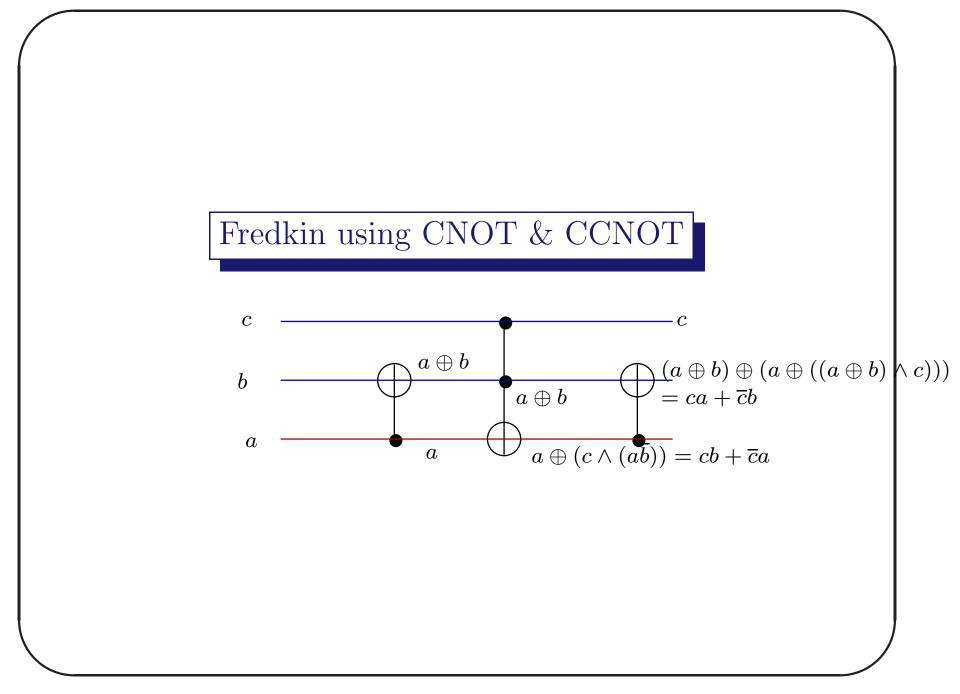


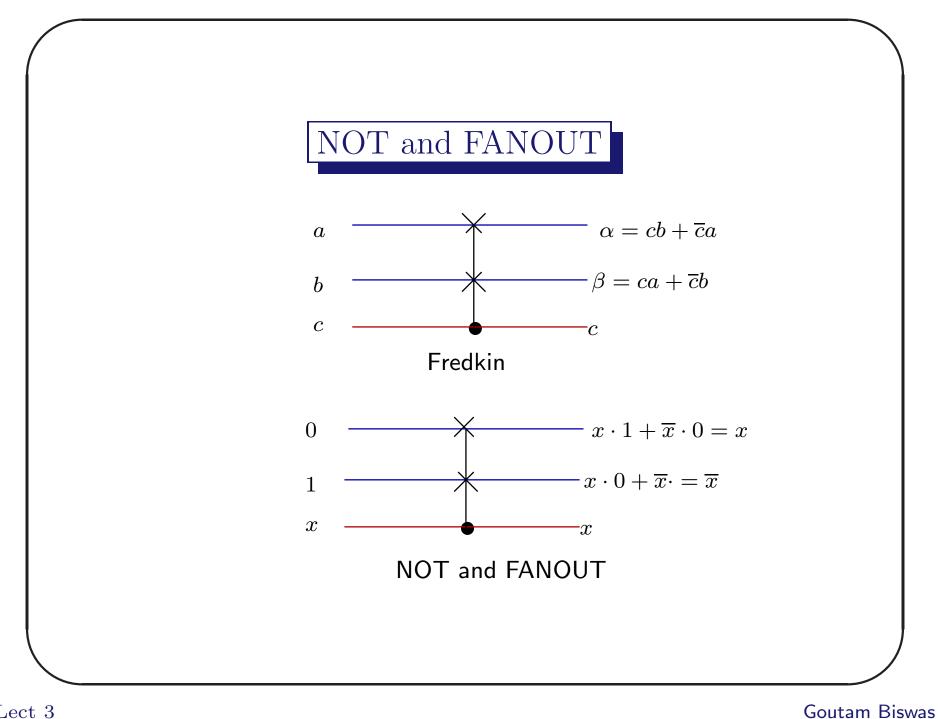
Clearly Fredkin gate is its own inverse.

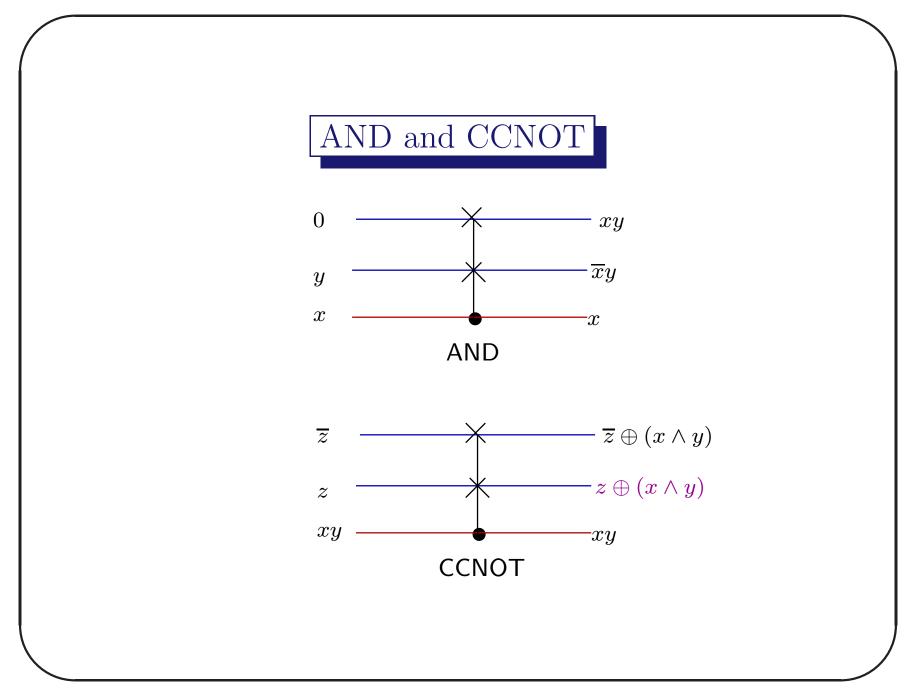


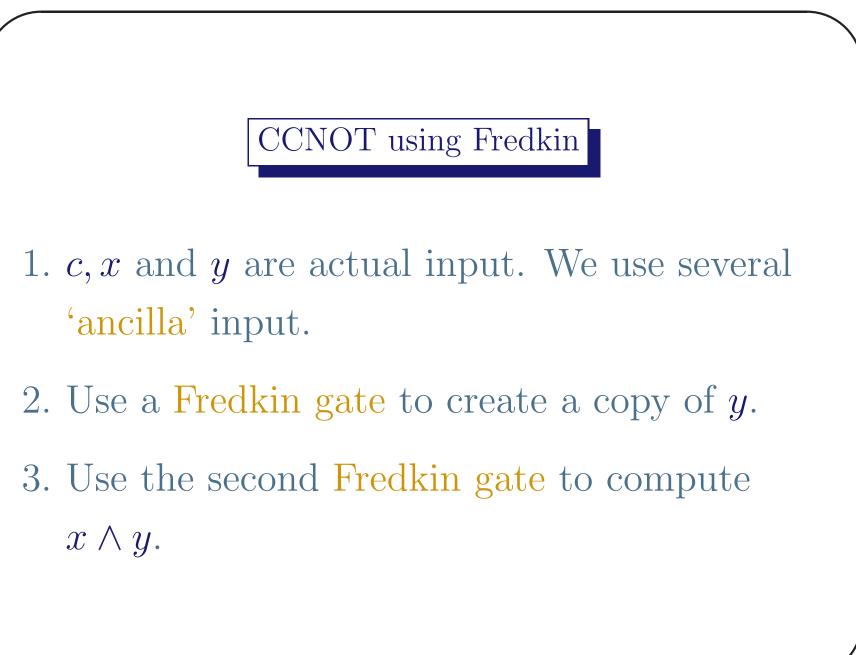


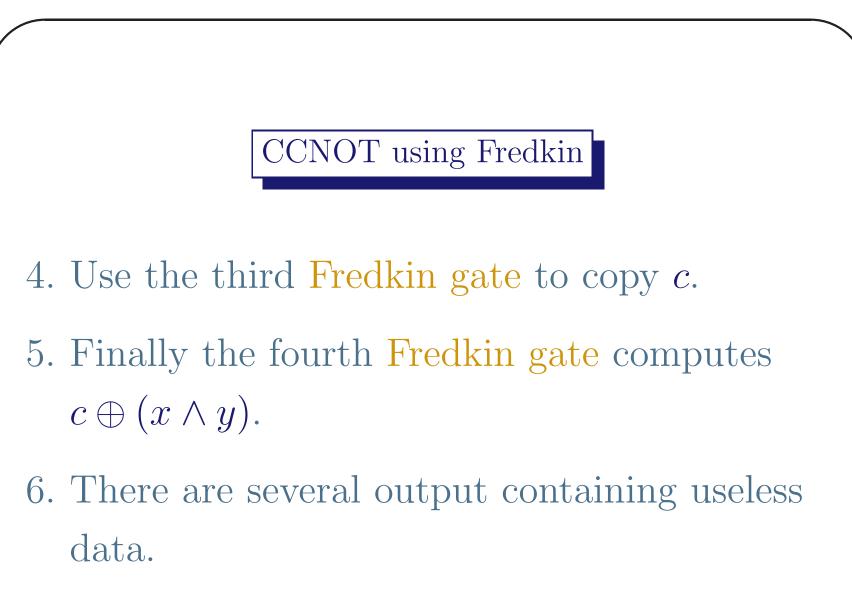
Both in case of Toffoli gate and in Fredkin gate we need some 'ancilla' input bits in state '0' or '1' to make them universal. They also produce useless outputs not used in subsequent stages.

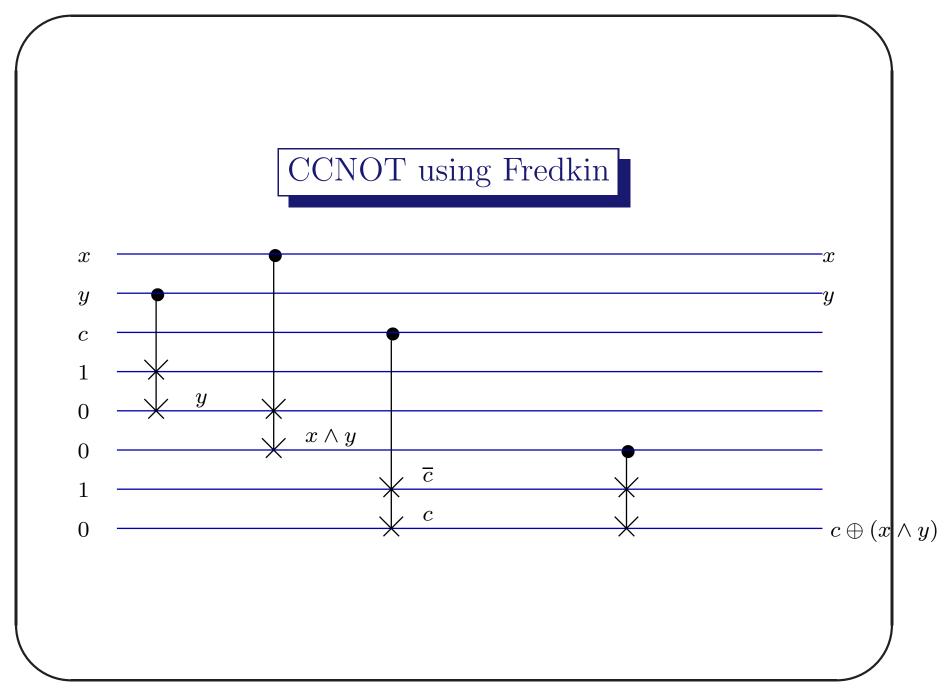


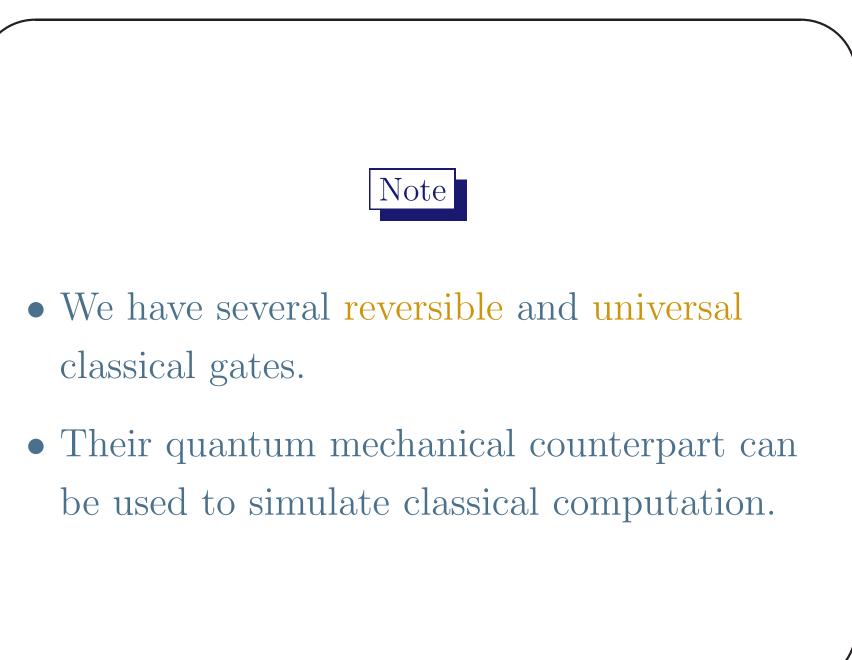












### Probabilistic Circuit

- The linear algebra formalism of classical bits and gates can be generalised to probabilistic circuits.
- Suppose a single-bit is at state 0 with probability p and at state 1 with a probability 1 p.
- It is represented as a 2-dimensional real vector  $\begin{bmatrix} p \\ 1-p \end{bmatrix}$ , where  $p \in [0, 1]$ .

## 1-bit Probabilistic Circuit

Applying our old transformation matrices we get,

$$c_{0} \begin{bmatrix} p \\ 1-p \end{bmatrix} = \begin{bmatrix} 1 & 1 \\ 0 & 0 \end{bmatrix} \begin{bmatrix} p \\ 1-p \end{bmatrix} = \begin{bmatrix} 1 \\ 0 \end{bmatrix},$$
$$\neg \begin{bmatrix} p \\ 1-p \end{bmatrix} = \begin{bmatrix} 0 & 1 \\ 1 & 0 \end{bmatrix} \begin{bmatrix} p \\ 1-p \end{bmatrix} = \begin{bmatrix} 1-p \\ p \end{bmatrix}$$
Transformation matrix for probabilistic bits is left stochastic (each column adds to 1).

### 2-bit Probabilistic Circuit

If we consider two bits so that

- the probability is  $p_0$  for the  $0^{th}$ -bit to be 0and the probability is  $p_1$  for the  $1^{st}$ -bit to be 0.
- The probabilities of 00, 01, 10, 11 are  $p_1p_0, p_1(1-p_0), (1-p_1)q_0, (1-p_1)(1-p_0)$ respectively.

#### 2-bit Probabilistic Circuit

The joint probabilities of two bits may be represented as the following tensor product -

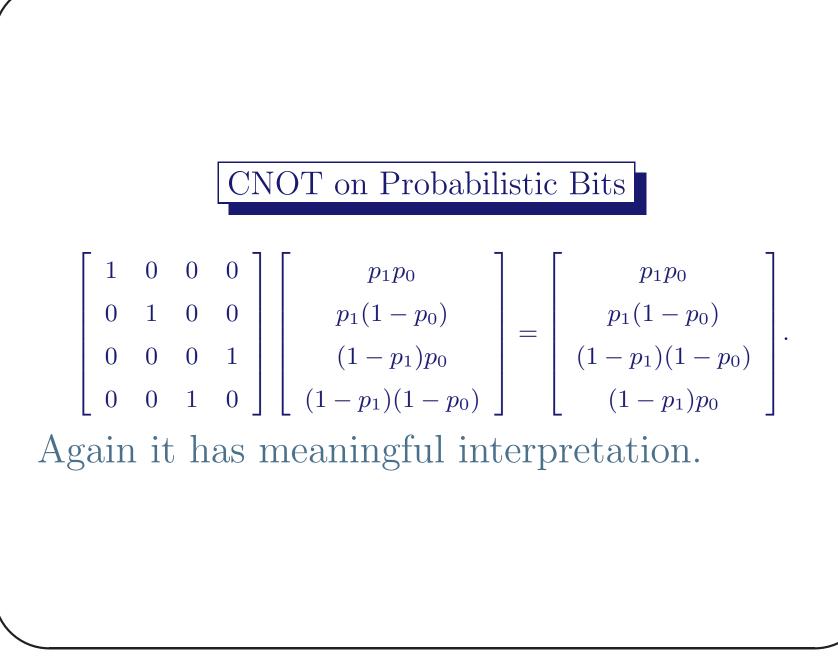
$$\begin{bmatrix} p_1 \\ 1-p_1 \end{bmatrix} \otimes \begin{bmatrix} p_0 \\ 1-p_0 \end{bmatrix} = \begin{bmatrix} p_1p_0 \\ p_1(1-p_0) \\ (1-p_1)p_0 \\ (1-p_1)(1-p_0) \end{bmatrix}$$

#### Xor-And on Probabilistic Bits

If we apply Xor-And transformation on two probabilistic bits, we get

$$\begin{bmatrix} 1 & 0 & 0 & 0 \\ 0 & 0 & 0 & 1 \\ 0 & 1 & 1 & 0 \\ 0 & 0 & 0 & 0 \end{bmatrix} \begin{bmatrix} p_1 p_0 \\ p_1 (1 - p_0) \\ (1 - p_1) p_0 \\ (1 - p_1) (1 - p_0) \end{bmatrix} = \begin{bmatrix} p_1 p_0 \\ (1 - p_1) (1 - p_0) \\ p_1 (1 - p_0) + (1 - p_1) p_0 \\ 0 \end{bmatrix}$$

The interpretation of transformation makes sense.



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- [MNIC] Quantum Computation and Quantum Information by Michael A Nielsen & Isaac L Chuang, Pub. Cambridge University Press, 2002, ISBN 81-7596-092-2.
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- [AAM] Classical and Quantum Computation by A Yu Kitaev, A H Shen & M N Vyalyi, Pub. American Mathematical Society (GSM vol 47) 2002, ISBN 978-1-4704-0927-2.

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